1 Geophysical Trends inferred from 20 years of AIRS **infrared global observations**

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Key Points:

or assimilation into Numerical Weather Prediction Models.

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¹⁴ Abstract

 NASA's Atmospheric Infrared Sounder has been in near-continuous operation since 16 September 2002. The \sim 3 million daily spectral observations contain detailed informa- tion about surface and atmospheric temperature, water vapor and trace gases such as $CO₂$ and CH₄, as well as clouds and aerosols. In this paper we obtain climate thermo- dynamic trends using 20 years of AIRS observations by working exclusively with the trends observed in the AIRS radiance time series. This is achieved by first binning the observed ²¹ spectra into nominal 3×5 degree latitude/longitude spatial subsets using 16 day inter- vals, after which a quantile-based algorithm selects nominally clear scenes for each grid box in order to construct the clear scene radiance spectrum time series. De-seasonalized spectral anomalies and spectral trends are then obtained from the time series, which are converted into geophysical trends using a physical retrieval for each grid box. This approach is completely different from traditional operational use of infrared data for trending, whereby are—anomalies/trends are generated either after daily retrievals, or after assimilation into NWP ²⁸ models. Our approach rigorously ties the derived geophysical trends to the observed ra- diance trends, and requires orders-of-magnitude fewer computational resources and time than re-analysis or traditional Level 2 retrievals. The retrieved trends are compared to trends derived from four other products : ERA5, MERRA2 reanalysis model fields and the NASA Level3 AIRS v7 and NASA Level 3 CLIMCAPS v2. Our retrieved surface tem- perature trends agree quite well with ERA5 re-analysis, CLIMCAPS L3 and the GISS surface climatology trends. Atmospheric temperature profile trends exhibit some vari- ability amongst all these data sets, especially in the polar stratosphere. Water vapor pro- file trends are nominally similar amongst all data sets except for the AIRS v7 which ex-37 hibits trends with a different sign in the mid troposphere. Note that infrared sounders lose water vapor sensitivity close to the surface making intercomparisons of column water ³⁹ trends problematic. Spectral closure between observation trends versus those computed by running all the NWP re-analysis and official NASA L3 monthly fields though a (clear sky) radiative transfer code is discussed, with the major differences arising in the wa-ter vapor sounding region.

43 Plain Language Summary

44 December 2021 and a metal sounders, designed for weather forecasting purposes, as have been in orbit around the Earth for a long enough time to enable anomaly and trending ϵ studies for climate purposes. Traditionally their daily obtained radiance data has been a used for operational atmospheric state retrievals, or assimilation into Numerical Weather as Prediction models, after which climate anomaly studies are made. In this paper we use the raw radiance spectral data to form radiance anomalies and trends, after which we so do a one step atmospheric state retrieval. This novel approach has the benefit of using a only stable channels together with easily understood assumptions and well tested retrieval algorithms to do the trend or anomaly geophysical retrieval, which has full error characterization. 53

⁵⁴ 1 Introduction

 NASA's Atmospheric Infrared Sounder (AIRS) became operational in September 2002, as the first of the new generation of low noise, high stability hyperspectral sounders, making \underline{Top} of Atmosphere (TOA) radiance observations at a typical 15km (at nadir) horizontal resolution. Follow on instruments with similar characteristics and abilities in- clude Eumetsat's Infrared Atmospheric Sounding Interferometer (IASI) and NOAA's Cross Track Infrared Sounder (CrIS), operational since June 2007 and March 2012 respectively. The latter two already have follow on missions planned till the 2040s, and together these ϵ_2 three sounders will provide scientists with a 40 year high quality, near continuous ob-servational dataset for climate anomaly and trending studies.

⁶⁴ Infrared radiances contain a wealth of information, including but not limited to sur-⁶⁵ face temperature, atmospheric temperature and water amount, and mixing ratios of green- ϵ house gases such as carbon dioxide CO₂, CH₄ and N₂O. <u>Measurements by visible imagers</u> **or** which have ∼1 km horizontal resolution or better King et al. (2013) suggest global cloud os free fractions of ∼ 30%, but the 15 km footprint of typical sounders means at most 5% **ω**of the hyperspectral observations can be considered "cloud-free." Current operational NASA ⁷⁰ L2 products come from cloud clearing the observed radiances, which introduces errors rand use the method of cloud clearing on observed radiances in partly cloudy scene conditions retrieval. The cloud clearing method solves for an estimate ✿✿ ra of clear column radiances by examining adjacent Fields of View (FOVs) to estimate the ra cloud effects on observed allsky radiances, assuming any differences are solely due to different retrieval amounts in each FOV, and significantly increases geophysical retrieval yields (to ⁷⁶ about 50-60%) Smith and Barnet (2023). This does introduce increased noise in the cloud σ cleared radiances of the lower atmosphere sounding channels; in addition the subsequent τ_8 retrieval depends on the first guess (which is a neural net for AIRS v7 and MERRA2 re-⁷⁹ analysis for CLIMCAPS v2). The The reader is referred to Susskind et al. (2003); Smith and Barnet (2020) ✿✿ ⁸⁰ Susskind et al. (2003); Smith and Barnet (2020, 2023) for more details. ⁸¹ In this paper we work directly in radiance space and form either anomalies or trends

⁸² from the underlying well characterized and understood radiances Strow and DeSouza-83 Machado (2020), in order to do a geophysical trend or anomaly retrieval. The work pre-⁸⁴ sented here, once the averaged/sorted data is available, can be processed in hours to days, ⁸⁵ and can be duplicated by small research groups with ease. Moreover, our novel approach ⁸⁶ has zero temperature *a-priori* and minimal water vapor *a-priori*. This completely sidesteps $\frac{1}{87}$ time variability and the accuracy of the *a-priori* which causes errors in the retrievals, ⁸⁸ and ensures our work examines trends directly inferred from the radiances versus those from traditional methods, leading. This leads to more unbiased results that directly high-⁹⁰ light the conditions (for example stratospheric water vapor) where the sensor has lim-⁹¹ ited sensitivity.

 The approaches used in this work are therefore very different than climate anoma- lies or trends from reanalysis products or traditional Level 2 retrievals, neither of which ₉₄ are tailored for climate trends. Reanalysis uses a wide range of observations and are only created within very large organizations, and represent the most commonly used climate ⁹⁶ data sets. They products assimilate individual sensor scenes from many different instruments, and may have discontinuities as different instruments come online or go offline. Tradi- tional Level 2 (and Level 3 products derived from Level 2) retrieve the atmospheric state for individual scenes (or effective cloud-cleared radiance derived from a 3x3 grid of in- dividual scenes). Both reanalysis and Level 2 products require large computational re- sources, that preclude full dataset re-processing to help fully understand trends. A main characteristic of traditional L2 retrievals is the requirement for a good *a-priori* state for each inversion, making errors in the *a-priori* difficult to distinguish from true variabil-ity in the data, especially with regard to trends.

 The stability and accuracy of the AIRS instrument is documented in recent work on analyzing 16 years of AIRS radiance anomalies over cloud-free ocean Strow and DeSouza- Machado (2020). Geophysical retrievals on the anomalies yielded $CO₂$, CH₄, N₂O and surface temperature time series that compared well against in-situ data from NOAA Global Monitoring Laboratories (GML) trace gas measurements and NOAA Goddard Institute of Space Studies (GISS) surface temperature data respectively. A significant difference between this paper and Strow and DeSouza-Machado (2020) is the nominally clear scenes used in this paper are selected uniformly from all over the Earth, while the clear scenes in the latter were zonal averages which were sometimes concentrated in certain regions.

¹¹⁴ In this paper we expand upon our initial zonal clear sky analysis, to derive geophys-¹¹⁵ ical trends from 20 years (September 2002 - August 2022) of AIRS measurements over \sim 3 \times 5 degree tiles covering the Earth, chosen such that the number of observations

 in each tile is roughly equal. An important concept introduced is spectral closure, whereby the observed clear sky spectral radiance trends are compared to spectral trends produced by running the monthly reanalysis or official NASA retrieved AIRS L3 products through an accurate clear sky radiative transfer code; close agreement in different sounding re g ₁₂₁ gions (such as 640-800 cm⁻¹ for temperature and CO₂, 1350-1640 cm⁻¹ for water vapor, 1000-1150 cm⁻¹ for O_3) between the computed and actual observed spectral trends imply that trends from those geophysical parameters used in the computations are re- alistic while disagreement suggests otherwise. A companion paper will utilize the geo- physical trend results to derive Outgoing Longwave Radiation (OLR) trends and non- local clearsky feedback parameters. Nominally clear scenes for each tile are picked out using a quantile approach; from the time series, radiances trends are made over the en-¹²⁸ tire Earth, from which geophysical trends are retrieved.

 Observed infrared spectral trends from AIRS has already been a focus of earlier work by Huang et al. (2023) who studied a slightly shorter time period (2002-2020) while Raghuraman et al. (2023) converted the radiances to Outgoing Longwave radiation (OLR), but neither study involve retrievals from spectral trends to geophysical trends. Instead they convert various model trends (such as ERA5) to spectral trends and compare against the observed spectral trends. Our earlier work shows we can accurately account for the effects of GHG forcings (Strow et al., 2021). In this paper we remove these GHG forc- ings from the observed AIRS spectral trends to concentrate on atmospheric temperature and water vapor and surface temperature, while the papers by Huang et al. (2023); Raghu- raman et al. (2023) include the GHG forcings in the model generated spectral trends. Another noteworthy examination of the time evolution of high spectral resolution infrared radiances (converted to spectral outgoing longwave radiation (OLR) fluxes) by (Whitburn et al., 2021) covered 10 years (2007-2017) of IASI observations. They confirmed that the IASI-derived fluxes agreed well with increases in GHG gas concentrations and El-Nino Southern Oscillation (ENSO) events within that time frame. A more recent paper Roemer et al. (2023) used the 10 year IASI data to derive allsky longwave feedback spectral com-145 ponents (water vapor, CO_2 , window, ozone) and total values, while also estimating clear sky feedback values.

 We will refer to our results as the AIRS Radiance Trends (AIRS_RT). Compar- isons are made against monthly output from the European Center for Medium Weather Forecast fifth generation reanalysis (ERA5) Hersbach et al. (2020) and NASA's second generation Modern-Era Retrospective analysis for Research and Applications (MERRA2) Gelaro and Coauthors (2017), and also against the official monthly AIRS L3 products which are AIRS v7 L3 Susskind et al. (2014); Tian et al. (2020) and CLIMCAPS v2 L3 Smith and Barnet (2019, 2020). Detailed geophysical trends and spectral closure stud- $_{154}$ ies are presented for the ascending (daytime (D)), descending (nightime (N)) and D/N averages.

2 Datasets used in this study

 Three main types of datasets are used in this study. The first is the AIRS L1C ra- diance dataset we analyzed for this paper, which has both daytime (D) and nightime (N) (ascending and descending) views of the planet. Second is the monthly operational L3 retrieval data, which are the AIRS v7 and the CLIMCAPS v2 products, also separated $_{161}$ into D/N data. Finally we also compared to trends from ERA5 and MERRA2 monthly reanalysis model fields. The ERA5 monthly dataset is available in 8 averaged time steps, 163 so we match to the average AIRS overpass times and compute (D/N) data over the 20 years, while MERRA2 monthly model fields are only available as one time step; included here for completeness we mention the NASA GISS surface temperature dataset, which ¹⁶⁶ like MERRA2 is only available as one set per montha monthly mean. This means four of the datasets : AIRS_RT (from AIRS L1C), AIRS L3 and CLIMCAPS L3, and ERA5 168 are separable into D/N , while the other two (MERRA2 and GISS) are only available as

¹⁶⁹ a diurnal averaged value. We describe these datasets in more detail below. In addition ¹⁷⁰ we also briefly mention other datasets that we use.

¹⁷¹ 2.1 The AIRS instrument and L1C dataset

 The Atmospheric Infrared Sounder (AIRS) on board NASA's polar orbiting EOS/Aqua platform has 2378 channels, covering the Thermal Infrared (TIR) spectral range (roughly $_{174}$ 649-1613 cm⁻¹) and shortwave infrared (2181-2665 cm⁻¹). The full widths at half maximum satisfy $\nu/\delta\nu \sim 1200$. The (spectral dependent) noise is typically $\leq 0.2K$. The orig- inal L1b radiance dataset suffers from spectral gaps and noise contaminated data as de- tectors slowly fail. These limitations are addressed using a 2645 L1c channel dataset, where spectral gaps and some of the noise "pops" are filled in using principal component recon- struction Manning et al. (2020) and is the dataset used to subset radiances analyzed in this paper. However we note that the results described in this paper used only the ac- tual observed radiances in pristine, stable channels described in Strow et al. (2021) and none of the synthetic channels. The Aqua platform is a polar orbiting satellite with 1.30 am descending (night time over equator) and 1.30 pm ascending (daytime over equator) tracks. Each orbit takes about 90 minutes, with the 16 passes yielding almost twice daily 185 coverage of the entire planet. About \sim 3 million AIRS spectral observations have been obtained daily since AIRS became operational in late August 2002. The instrument has provided data almost continuously since then though there have been some shutdowns (each spanning a few days) such as during solar flare events.

 In this paper we use the re-calibrated 2645 channel L1C radiance data Strow and DeSouza-Machado (2020) instead of the 2378 L1B data. 20 years (spanning September 1, 2002-August 31, 2022) of AIRS L1C radiance data is gridded into 4608 tiles covering the Earth : 72 longitude boxes which are all 5° in width, and 64 latitude boxes which are 193 approximately 2.5° in width at the tropics but wider at the poles to keep the number of observations per 16 day intervals (which is the repeat cycle of the AIRS orbit on the Aqua 195 satellite) roughly the same. This way there are \sim 12000 observations per 16 days per tile. 196 which are roughly equally divided between the ascending/daytime (D) and descending/nigh-¹⁹⁷ time (N) tracks. In this paper we discuss results for both the ascending and descending tracks using a retrieval based on the longwave (LW) and midwave (MW) regions of the spectrum (640-1620 cm⁻¹ or 6-15 μ m).

 In this paper our trend retrievals use only the AIRS channels are stable in time, as quantified in Strow et al. (2021). For example the shortwave (SW) channels are drift- $_{202}$ ing at a higher rate than the LW/MW channels, which can lead to incorrect surface tem- perature rates, and are avoided in this paper. Similarly there are are many channels in the LW and MW whose detectors are drifting in time, and which are also not used here. For example there are some higher wavenumber (shorter wavelength) channels past the ozone band which have a significant drift in time, possibly due to changes in the polar- ization of the scan mirror coating with time. Therefore compared to other AIRS oper- ational products used in this paper, our results use channels that are demonstrated to have high stability Strow et al. (2021). We do note that some of the observed drifts in the AIRS channels stabilized after 6 years, so their impact is reduced when looking at 20 year trends.

 In "clear sky" scenes, the window region would be dominated by the effects of water vapor continuum absorption, which is largest for hot and humid tropical scenes and almost ₂₁₄ negligible for cold, dry scenes. Scattering and absorption by liquid and ice clouds also $_{215}$ affects the window region (800-1000 and 1100-1250, and 2400-2800). For each tile, we 216 use the 1231.3 observation as our window channel (AIRS L1C channel ID = 1520), and ²¹⁷ form the quantiles of the observed Brightness Temperature (BT) for each 16 day observation period. BT 1231 therefore serves as a measure for the cloudiness of an observation : if there are no or low clouds, it will effectively measure the surface temperature, but as the

 clouds get thicker and higher, it will measure colder temperatures. Quantiles 0.50, 0.80, 0.90, 0.95 and 0.97 were among those chosen; the first would be considered the "median" observation, containing clear and cloudy scenes. In a subsequent section we show Q0.90 onwards can be considered "almost free of clouds." Our retrievals using this dataset are ²²⁴ referred to as in what follows.

2.2 Reanalysis Model fields

 The ERA5 fifth generation reanalysis product from the European Center for Medium Range Weather Forecasts is freely available on monthly timescales from the Copernicus $\text{Climate Data Store. This monthly dataset is output at 37 pressure levels at } 0.25^{\circ} \text{horizontal}$ resolution Hersbach et al. (2020), which is further subdivided into eight 3-hour averages per month (corresponding to 00,03,06,...21 UTC). For each month from September 2002- August 2022 we downloaded the surface temperature and pressure fields, as well as at- mospheric temperature, water vapor and ozone fields. These are then colocated to each tile center using 2D spatial interpolation, as well as time interpolated according to the average AIRS overpass time as a function of month. From the resulting monthly time- series of reanalysis model fields for each tile, we generated (a) thermodynamic trends for surface temperature, air temperature, water vapor and ozone model fields (b) a 20 year average thermodynamic profile in order to produce jacobians for the linear trend retrievals (c) by using the model fields as input to the clear sky SARTA radiative transfer code Strow, Hannon, DeSouza-Machado, et al. (2003) a monthly time series of clear sky radiances for each tile was generated, from which we could compute radiance trends. We did this ²⁴¹ for both the ascending and descending datasets.

 The MERRA version 2 (MERRA2) re-analysis used in this paper is the second gen- eration Gelaro and Coauthors (2017) product from NASA's Global Modeling and As- similation Office. The monthly data we use is available on 42 pressure levels at a hor- $_{245}$ izontal resolution of $0.5° \times 0.625°$, but only one monthly mean diurnally averaged out- put is available per month. Similar to the ERA5 output, we colocated the MERRA2 sur- face temperature, atmospheric temperature, water vapor and ozone fields to our tile cen- ters for each month starting September 2002 in order to produce a time series of radi- ance and model output, from which radiance and thermodynamic trends could be com- puted for comparisons against other datasets in this study; similar to above we also gen- erated a monthly time series of clear sky radiances for each tile, from which we could com-pute clear sky radiance trends based on MERRA2.

 The NASA Goddard Institute of Space Studies (GISS) surface temperature data v4 surface temperature data 2023 (2005); Lenssen et al. (2019) is a monthly dataset based primarily on near surface temperatures land stations, and data from ships and buoys. As with MERRA2 we obtained one monthly mean dataset per month, which we could not separate into descending (N) or ascending (D) tracks.

2.3 AIRS L3 Products

 NASA routinely produces two retrievals from the AIRS L1C data observed each day, which are AIRS v7 Susskind et al. (2014); Tian et al. (2020) and CLIMCAPS v2 Smith and Barnet (2019, 2020). Both use the cloud clearing process but there are sig- nificant algorithmic differences; in particular the AIRS v7 product is initialized by a neu- ral net, while CLIMCAPS uses MERRA2 for its initialization. The L2 products are then individually turned into L3 monthly products, for both the ascending (daytime) and de- scending (nighttime) data. The timeseries of thermodynamic profiles were used as in- put to the clear sky SARTA RTA to generate radiances, after which radiance trends and ²⁶⁷ thermodynamic trends are also produced.

²⁶⁸ 2.4 Other L3 Products

²⁶⁹ The Microwave Limb Sounder (MLS) monthly binned water vapor (H2O) mixing ratio dataset Lambert et al. (2021), which contains data at spatial coverage $\pm 82^\circ$ lat-²⁷¹ itude, at a spatial resolution of $4° \times 5°$ and useful vertical range between 316 and 0.00215 ²⁷² hPa was used in this paper to improve retrieval trends in the upper atmosphere.

²⁷³ 3 Filtering the Observational Data for clear scenes

 Here we discuss the "clear scene" selection from all the observed data stored for each ₂₇₅ of the 72×64 tiles. Ideally we would prefer to use a MODIS cloud fraction product (1) km) colocated to the 15 km AIRS footprints, but this is presently unavailable. Our ear- lier work used an uniform clear flag over ocean Strow et al. (2021) which will not work well over land because of surface inhomogeneity. In this section we discuss an alterna- tive clear filter based on the hottest 10 percent of AIRS observations that are present inside any 16 day tile, over any location.

²⁸¹ 3.1 Observed BT1231 Distributions

The left panel of radiances measured in thermal infrared window region (800-1000 cm⁻¹ and 1100-1250 cm⁻¹) are dominated by the effects of the surface temperature, water vapor continuum absorption and cloud/aerosol effects. The effects of water vapor continuum absorption is largest in hot and humid tropical scenes (depressing the observations relative to surface temperatures by about $5-6$ K, which reduces to about 2 K at \pm 50°) and is almost negligible for cold, dry scenes (less than 1 K). Scattering and absorption by liquid and ice clouds also affects the window region (Deep Convective Clouds can depress the ²⁸⁹ window channel observations by as much as 100 K relative to surface temperatures). For each tile, we use the 1231.3 cm^{-1} observation as our representative window channel $(AlRS)$ 291 L1C channel ID = 1520), as it is minimally impacted by weak water vapor lines. Changed to Brightness temperature (BT) the observation in this 1231.3 cm^{−1}channel (BT1231) therefore serves as a measure for the cloudiness of an observation : if there are no or low or optically thin clouds, it will effectively measure the surface temperature, but as the clouds get thicker and higher, it will measure the cold cloud top temperatures. For any ²⁹⁶ tile during any 16 day observation periods, we can compute quantiles $\mathcal Q$ based on the est observed BT1231 to screen between cloudy and partially clear scenes. We chose different es BT1231 quantiles (so quantile Q0.XY will have a numerical value $BT1231_{Q0,XY}$ associated ess with it) and show below the data contained between $Q0.90$ and $Q1.00$ can be considered "almost free of clouds." 300

 $F_{\text{figure 1 shows the zonally averaged histograms for a all the BT1231 observations}}$ of a chosen 16 day timestep in the form of a zonally averaged histogram (normalized as probability distribution functions (PDFs)), with latitude on the vertical axis and BT1231 and on the horizontal axis. The colorbar is the PDF value, and we used data spanning Au-³⁰⁵ gust 27, 2012 - September 11, 2012 which is approximately half way through the 20 year ³⁰⁶ AIRS mission dataset used in this paper. The colorbar is the mean histogram (normalized ³⁰⁷ probability distribution functions (PDFs)) using the data in that 16 day time period. From this we plot the curves show the zonally averaged BT1231 values of the minimum (Q0.00) 309 Q0.00) in dark cyan, mean (thick red), mean, median (Q0.50Q0.50 in orange), maximum $_{310}$ (Q1.00Q1.00 in light cyan); also shown are the a handful of other zonally averaged BT1231 v_{311} values of Q0.80, Q0.90 , for example $Q_{0.80}$, $Q_{0.90}$ (thick black curve), Q0.95 and Q0.97. ³¹² The BT1231 channel has the lowest expected absorption due to water vapor in the longwave 313 portion of the spectrum, and so is expected to sense the surface temperature unless the 314 seene is cloudy in which case it would be expected to sense the cloud top temperature. ³¹⁵ In this way the histogram should exhibit the characteristics of the cloud conditions observed in the 16 day period. $Q0.95$ and $Q0.97$. The distributions are skewed to the left (negative ar skewness), as confirmed by the mean being less than the median. We also point out that

even $Q0.80$ sees much of the surface from the southern tropics to the northern polar region. The 220 K horoizontal axis cutoff means we do not see the very cold (190 K) observations over the winter Antarctic. 320

 $\frac{1}{221}$ $\frac{1}{21}$ figure shows the expected qualitative features, for example (1) the tropical PDFs ³²² peak at around 295 K, but show some warmer observations, as well much colder obser-³²³ vations (below 230 K) corresponding to Deep Convective Clouds (DCC); this gives a dy- $_{324}$ namic range of almost 100 K at the tropics (2) the BT1231 observed over the Southern ³²⁵ Polar (polar winter) regions are much colder than the BT1231 observed over the Northern Polar (polar summer) regions and (3) the reddish peaks in the $30°N - 40°N$ are a com-³²⁷ bination of the marine boundary layer (MBL) clouds and warmer summer land temper-³²⁸ atures.

³²⁹ It is evident the distributions are skewed to the left (negative skewness), as confirmed ³³⁰ by the mean being less than the median. We also point out that even Q0.80 sees much ³³¹ of the surface from the southern tropics to the northern polar region. The right panel 332 of Figure 1 shows the same information, except presented as a cumulative histogram, with α value of 0 at the hot end (340 K) and 1 at the cold end (180 K); again one sees the Q0.90 $\frac{1}{334}$ quantile envelopes the hottest 10% of the observations as expected. The cutoff of 220 $\frac{1}{335}$ K in the plots does not allow the plot to extend to show the very cold (190 K) observations ³³⁶ over the winter Antarctic.

 337 Zonally averaged BT1231 histograms for an $2012/08/27 - 2012/09/11$ timespan (colorbar) 338 and quantiles (curves). The thick black curve is the Q0.90 quantile (and above) used in ³³⁹ this paper, and is very close to the maximum. The left hand panel shows the normalized ³⁴⁰ histogram (probability distribution function) as a function of latitude and temperature ³⁴¹ bin; the right hand panel shows the cumulative distribution function, though starting $_{342}$ from the hotter side (cdf(340 K) = 0.0, cdf(180 K) = 1.0).

³⁴³ Measurements by visible imagers which have ∼ 1 km horizontal resolution or better ³⁴⁴ King et al. (2013) suggest global cloud free fractions of ∼ 30%, but the 15 km footprint ³⁴⁵ of typical sounders means at most 5% of the hyperspectral observations can be considered ³⁴⁶ "eloud-free." In the tropics, the higher amounts of water vapor means the observed BT1231 347 for a clear scene would be reduced by a $5-6$ K due to water vapor continuum (which on 348 average reduces to about 2 K at ± 50 , and 1 K at the polar regions). Figure 1 shows on average the on average the cloud effect at the tropics is an additional modest 20 K (difference between $\sqrt{Q0.90}$ and $\sqrt{Q0.90}$ and $\sqrt{Q0.50}$ compared to the 100 K dynamic range. ³⁵¹ This is because the cloud fractions and cloud decks in the individual observations have ³⁵² effectively more clouds (with larger cloud fraction in the FOV) lower in the atmosphere ³⁵³ than higher up; the net effect is that in the window region the atmosphere is on average radiating from the lower (warmer) altitudes, and so $Q0.80$ to $Q1.00$ onwards spectra $\frac{\text{whose BT1231 values are larger than BT1231}_{\text{Q}_0,80}$, see much of the surface emission as ³⁵⁶ well.

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We now use the above plots to select "almost clear" scenes. For any one tile, we define $\text{set }\Psi_{0,XY}$ to have all observations i whose BT1231 lies between quantiles $\mathcal{Q}0.XY$ and 360 $Q1.00, \{i\mid BT1231_{Q0,XX} \leq BT1231(i) \leq BT1231_{Q1.00}\}$. In what follows in this subsection ³⁶¹ we use an "integrated" or "cumulative" quantile wherein Q0.90 now means all scenes between ₃₆₂ Q0.90 and Q1.00 (maximum observed BT) are considered. XY is the radiances averaged over all the observations i which are in the set $\Psi_{0,XY}$, namely 363

$$
r_{Q0,XY}(\nu) = \frac{1}{N_{0,XY}} \sum_{i \in \Psi_{0,XY}} r_i(\nu)
$$
 (1)

Figure 1. Zonally averaged BT1231 normalized histograms (probability distribution functions) as a function of latitude and temperature bin, for an 2012/08/27 - 2012/09/11 timespan (colorbar) and quantiles (curves). The thick black curve is the Q0.90 quantile (and above) used in this paper, and is very close to the maximum Q1.00 quantile.

 $\frac{1}{2}$ where $r_i(\nu)$ are the $N_{0,XX}$ individual observations in set $\Psi_{0,XX}$. In this section we only ass use the $\nu = 1231$ cm⁻¹channel, but in later sections we easily form averages for all 2645 asse channels, at any 16 day time step for any tile.

³⁶⁷ To further investigate if the scenes chosen using this definition can be considered $\frac{1}{2}$ We tested different quantile sets $\Psi_{0,XX}$ to see which one can reliably be considered to provide a nominally "cloud free" , we compare to global dataset, and chose the Q0.90 average areased over the $\Psi_{0.90}$ set) as the one to use for the rest of this paper, an unless explicitly stated otherwise. The tests primarily involved comparisons to scenes ³⁷² produced by the uniform/clear sky filter described in Strow and DeSouza-Machado (2020) ³⁷³ for the same August 27, 2012 - September 11, 2012 sixteen day timespan. This latter fil-³⁷⁴ ter selects clear scenes by both testing for uniformity (to within 0.5 K) across a 3×3 ³⁷⁵ grouping of AIRS scenes and also using a criteria that the observed window channel ob- 376 servations should be within ± 4 K of clear-sky simulations using thermodynamic param-³⁷⁷ eters supplied by reanlysis reanalysis models. The results are shown in the left hand plot 378 of Figure 2, plotted on a $1° \times 1°$ grid. We note in this plot the uniform/clear scenes that are plotted are limited to those over ocean, and for solar zenith less than 90 \degree (daytime), ³⁸⁰ which automatically filtered out many of the views over the (wintertime) Southern Po- α ³⁸¹ lar region. Immediately apparent are the gaps produced by the uniform/clear filter e.g. ³⁸² in the Tropical West Pacific or off the western coasts of continents where there are clouds. 383 The gaps can be changed by e.g. changing the 4K threshold to allow more or fewer scenes ³⁸⁴ through the filter.

³⁸⁵ The center plot shows the scenes selected by the integrated for all tiles, the daytime selected for the Q0.90 filter average for the same time period, on the same $1^{\circ} \times$ ³⁸⁷ ¹°grid. Compared to the left hand plot, the spatial coverage is almost complete, as the 388 Q0.90 filter average always has the hottest 10% of the observations, the spatial coverage is almost complete : gaps are only visible . At this 1° resolution, used for comparison with the uniform/clear grid filter described in the previous paragraph, gaps are seen in regions

Figure 2. Clear scenes for the same $2012/08/27 - 2012/09/11$ timespan selected by (left) and uniform/clear sky filter and (center) the $Q0.90$ integrated filter average described in this paper. The right hand plot shows the mean (over ocean) observed BT1231 as a function of latitude, for the two selections; the difference is about $0 K \pm 1 K$ in most region except in the southern midlatitudes where the integrated Q0.90 filter-average produced scenes that were about 1 K cooler on average.

³⁹¹ where there are for example mountains, or in the desert regionswhere other areas are even 392 warmer. We note that increasing the quantile threshold to 0.95 or 0.97 did not introduce the gaps seen in the left hand (uniform/clear) map for example the local topography means observations over mountains would be colder than the surrounding coastal or plain regions. ass This is not a concern since zooming back out to the coarser $3^\circ \times 5^\circ$ tile resolution, will include Q0.90 data for the quantile and trending analysis. 396

³⁹⁷ To compare the mean observations we filter away remove the over-land and over-³⁹⁸ polar region data from the center plot. The right hand plot shows the mean observed BT1231 from the 1° \times 1° grid from the uniform/clear sky filter as a function of latitude, ⁴⁰⁰ compared to the 1[°] × 1[°]grid from the integrated Q0.90 scenes. The difference between ⁴⁰¹ the uniform/clear versus integrated Q0.90 filter average is within about 0.25 K \pm 1 K ⁴⁰² across the southern tropics to the northern midlatitudes, though the bias rises to about 403 1 K by about -50°S. We consider this an acceptable difference, as we could tune the thresh- $_{404}$ olds for the uniform/clear filter to e.g. change the areal coverage and/or number of clear ⁴⁰⁵ scenes and hence comparisons to the Q0.90 scenes.

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⁴⁰⁷ We ran these tests for three The results presented in this section have been checked ⁴⁰⁸ for robustness, using other 16 day intervals in 2012, spanning the four seasons. The overall $_{409}$ global bias and standard deviation for the 1231 channel between the co-located 1 \times 1grids ⁴¹⁰ by the uniform/clear filter and by the Q0.90 filter stayed fairly uniform, typically about $\frac{0.25 \text{ K } \pm 1 \text{ K}}{25 \text{ K } \pm 1 \text{ K}}$. From the information presented in this section, we We conclude that for any 16 day timestep the integrated radiances used in the Q0.90 filter average (a) pro-⁴¹³ duces almost complete spatial coverage of the Earth, (b) selects scenes whose average ⁴¹⁴ BT1231 is very close to the average BT1231 from scenes selected using an uniform/clear ⁴¹⁵ filter (c) trends from that quantile typically differ by less than \pm 0.002 K \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow ⁴¹⁶ the other quantiles and (d) this selection produces spectral trends which compare well ⁴¹⁷ against those obtained from the quality assured binned AIRS CCR data record Manning 418 (2022). Together these imply the integrated Q0.90 average is an acceptable proxy for "clear scenes". For the remainder of the paper we drop the word "integrated" and therefore con-⁴²⁰ sider Q0.90 as consisting of nominally clear observations whose BT1231 lies between the ⁴²¹ 90th quantile and hottest observation. <u>Our retrievals using this $Q_{0.90} \rightarrow Q_{1.00}$ averaged</u> dataset (shortened to Q0.90) is referred to as AIRS_RTin what follows. 422

423 3.2 Observed trends from the Q0.90 Quantiles

 Having selected the Q0.90 observations, for each tile the average radiance per 16 day interval is computed. With two sixteen day periods not available (Aqua platform ₄₂₆ or AIRS shutdowns during e.g. solar flare events) this gives a total of 457 time steps over 20 years. Anomalies are formed from this time series, and then de-seasonalized to give the spectral radiance trends and error estimates Strow and DeSouza-Machado (2020) us- $_{429}$ ing Matlab *robustfit*:

$$
r_{\text{observations}}^{16 \text{ days}}(t) \sim r_{\text{fit}}(t) = r_o + a_1 t + \sum_{i=1}^{4} c_i \sin(n2\pi t + \phi_i)
$$
 (2)

430 with a_1 and its associated uncertainty, both converted to brightness temperature (BT), ψ_{431} being the trends in K $\overline{/y_{\text{F,Y}}^{-1}}$. Using sub-harmonics in the fit did not produce any no-⁴³² ticeable change in the AIRS_RT retrievals (described below).

⁴³³ The left panel of Figure 3 shows the descending orbit (nightime) 20 year (Septem-⁴³⁴ ber 2002- August 2022) global averaged spectral observations for the five quantiles men-⁴³⁵ tioned above. We note the spectra in most of the plots in this section are weighted by ⁴³⁶ the *cosine(latitude)* of the tiles, unless otherwise stated. In addition we only show the $640-1640$ cm⁻¹ region, and ignore the shortwave 2050-2750 cm⁻¹ region since the AIRS ⁴³⁸ SW channels are drifting relative to the LW Strow and DeSouza-Machado (2020). Spec-⁴³⁹ tral averages constructed from Figure 1 would have this same behavior, namely that in the window region the mean spectrum of data populating the warmer quantiles integrated $\frac{1}{441}$ are on the order of $Q1.00-(Q0.80, Q0.90, Q0.95, Q0.97)$ as defined in Equation 1 are on the order of ⁴⁴² a Kelvin apart, and have about half/quarter that difference in the optically thicker re-⁴⁴³ gions dominated by $H_2O \text{ and/or } CO_2$ absorption respectively.

⁴⁴⁴ The right hand panel of Figure 3 shows (top) the trends and (bottom) the 2σ trend uncertainties for these quantiles, in $\text{Kelvin}/\text{year}$, Kyr^{-1} . We emphasize that the top right ⁴⁴⁶ panel shows that the spectral trends for the quantiles lie almost on top of each other; the difference between the Q0.50 and other trends is at most about $+0.003 \text{ K} / \text{y} \text{y} \text{y}^{-1}$ 447 (out of a 0.02 K $\sqrt{y_{\text{Fyl}}^{-1}}$ signal) in the window region (and about +0.0045 K $\sqrt{y_{\text{Fyl}}^{-1}}$ 448 ⁴⁴⁹ in the troposphere temperature sounding channels), or less than 10%. Similarly the largest ⁴⁵⁰ trend uncertainty in the bottom panel is for Q0.50. This implies that clouds effects in ⁴⁵¹ the infrared do produce the largest variability (blue curve) but on average for the infrared are not changing much, so the $+0.022 \text{ K } / \text{year_yr}^{-1}$ window region trends are dominated ⁴⁵³ by surface temperatures changes and to a lesser extent by water vapor changes.

 454 The–TOA radiances in the 15 um (700-800 cm⁻¹) region is dominated by the–are ϵ ₄₅₅ impacted by two effects (a) the increased optical depths due to increasing atmospheric ⁴⁵⁶ CO₂ increases; the effects of increasing are to make the atmosphere emit at leads to atmospheric ✿✿✿✿✿✿✿ a✿ -0.06 K/year signal for the troposphere; hidden in there are the ,✿✿✿✿ and✿✿✿✿ (b) ✿✿✿ the✿✿✿✿✿✿✿✿✿✿✿ ⁴⁵⁸ atmospheric emission from higher altitudes/lower temperatures, leading to an almost resulting in almost temperature increases (again about $+0.02 \text{ K}/\text{year}$); also yr^{-1}). Also of interest is the t_{460} trends in the stratosphere (650-700 cm⁻¹) changes which consists of a stratospheric cool- μ_{461} ing signal (negative) and emission higher up due to increased CO₂; combining to give 462 a net zero effect over 20 years, also seen in Raghuraman et al. (2023). The H_2O signal $\frac{1}{463}$ is evident in the 1400-1625 cm⁻¹ region, and is negative; in other words, increasing tem- 464 peratures have led to increased atmospheric amounts of H_2O , and the water vapor feed-⁴⁶⁵ back has reduced the amount of outgoing flux in that region. By extension, this also hap-⁴⁶⁶ pens in the Far Infrared regions affected by water vapor; current sounders do not make $\frac{467}{467}$ direct measurement in the10-600 cm⁻¹ region so at present this can only be inferred; how-⁴⁶⁸ ever in the near future it is anticipated the Far Infrared Outgoing Radiation Understand-⁴⁶⁹ ing and Monitoring (FORUM) mission Palchetti et al. (2020) will provide data to fill in ⁴⁷⁰ this important gap in the future.

Figure 3. 20 year trends from different observation quantiles. The left hand panel shows the mean globally averaged BT trends observations from 20 years of AIRS data, for quantiles $50Q0.50,800.80,900.90.950.95,970.97$ as described in the text. The right hand panel shows (top) the globally averaged trends from for those years different quantiles and (bottom) the spectral uncertainty in the trends. The nightime (descending) trends are shown in these plots.

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⁴⁷² 3.3 Observed trend changes over 20 years

⁴⁷³ The left panel of Figure ?? shows the mean descending (nightime) orbit Q0.90 brightness ⁴⁷⁴ temperature spectrum, for four time periods, all commencing on September 1, 2002 - the ⁴⁷⁵ periods are for 5,10,15,20years of data and end on August 31, 2007, 2012, 2017, 2022respectively. ⁴⁷⁶ As expected the mean cosine averaged observed BT is slightly over 284 K through most ⁴⁷⁷ of the longwave window region. The right hand panel of the same Figure ?? shows the ⁴⁷⁸ trends for the four time periods in the top, while the bottom shows the uncertainties. ⁴⁷⁹ Averaging over the inter-annual variability affects the trends, with the shortest/longest ⁴⁸⁰ time periods (5/20 years) having the largest/smallest spectral uncertainty as one would ⁴⁸¹ expect as inter-annual variability slowly becomes less important in the trends.

⁴⁸² Changes in AIRS observations over time spans of 05,10,15,20 years all beginning ⁴⁸³ on September 1, 2002. The left hand panel shows the mean globally averaged 90th quantile ⁴⁸⁴ BT spectra for those time periods. The right hand panel shows (top) the trends from ⁴⁸⁵ those years and (bottom) the spectral uncertainty in the trends. The nightime (descending) 486 trends are shown in these plots.

⁴⁸⁷ 4 Spectral closure : comparisons between observed and simulated spec-⁴⁸⁸ tral trends

 Previous work Strow and DeSouza-Machado (2020) has demonstrated that the ra- $\frac{490}{490}$ diances from AIRS are climate quality, if one restricts the channel set to the ∼ 450 chan- nel set that is largely immune to nonphysical drifts Strow et al. (2021). In this section we describe a way to test the quality of the monthly thermodynamic output from reanal- ysis and/or L3 products which are all in geophysical space, against the AIRS L1C ob- servational data which is in radiance space. This is accomplished by geolocating the monthly (ERA5) surface temperature, air temperature, water vapor and ozone fields to tile centers as described in Section 2.2, which are then input and run through the **SARTA** fast model Strow, Hannon, DeSouza-Machado, et al. (2003), for the entire 20 years. Spec- tral radiance trends were then computed from these time series of (clear sky) spectral radiances. The conversion of L3 retrieval and NWP reanalysis trends to a radiance time

Figure 4. 20 year zonally averaged spectral brightness temperature trends (in K \rightarrow \rightarrow \rightarrow \rightarrow \rightarrow for (left) AIRS Q0.90 observations and (right) clear sky simulations using ERA5 monthly model fields. The center panel shows the AIRS Q0.90 spectral uncertainties. The ERA5 simulations included linear trends of $CO₂$, $CH₄$ and N₂O, while the O₃ trends in ERA5 are from the reanalysis itself.

⁵⁰⁰ series, provides a rigorous check of their accuracy against the observed AIRS L1C ra-⁵⁰¹ diance trends which are validated to be highly accurate.

⁵⁰² A good reviewer might ask about the noise introduced by secant angle varying in the 16 day period. Che

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⁵⁰⁴ The simulations included realistic column linearly-increasing-with time mixing ra- $_{505}$ tios for CO₂, CH₄ and N₂O for the ERA5 spectra, as well as land or ocean surface emis-506 sivity co-located to tile centers together with view angles of about 22°. From these the ⁵⁰⁷ ERA5 spectral trends were derived similarly to what was described above for the AIRS ⁵⁰⁸ observation spectral trends.

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Figure 4 shows the descending (night) zonally averaged results in K \rightarrow yeary r^{-1} , al- $_{511}$ lowing us to compare the Q0.90 nominally clear AIRS observed spectral trends, to those 512 simulated using monthly ERA5 fields (without clouds). The center panel shows the spectral trend uncertainties from the observations, also in K $\sqrt{\text{year}yr}^{-1}$. In the next section ⁵¹⁴ we derive geophysical trends from these (AIRS observed) spectral trends, and the sim-⁵¹⁵ ilarities/ differences in geophysical trends can be partially understood from the similar- $_{516}$ ities/differences in the spectral trends. For example, the H₂O sounding region (1350-1600) cm^{-1}) shows roughly similar (positive) trends in the tropics and mid-latitudes; there are some slight differences in the high altitude channels (1450-1550 cm⁻¹ region). The fol-⁵¹⁹ lowing sections shows that there are subtle differences in these trends, which manifest ⁵²⁰ as differences in tropospheric water vapor trends. Observations and simulations both have positive dBT/dt in the 800-960,1150-1250 cm⁻¹ region, indicating surface warming; how-⁵²² ever the ERA5 simulation show more warming in the southern polar regions than do the ⁵²³ AIRS observations. In particular note the mean warming in the tropics is less than that 524 in the mid-latitudes, and the polar regions show the largest overall change in brightness temperature in the window region. Large differences are seen in the 10 um (1000 cm^{-1}) 526 O₃ sounding region, which are not surprising since ozone assimilation is not a primary ⁵²⁷ goal of ECMWF assimilation; here we do not address these as we focus on the changes ⁵²⁸ to the moist thermodynamic state. The window region trends computed using the ERA5 529 model are more positive in the Southern Polar region. Conversely the 640-700 cm⁻¹ spec-⁵³⁰ tral region is positive, especially in the tropics; however the observations show a net cool- $\frac{1}{531}$ ing trend away from the tropics, compared to the ERA simulations. This demonstrates $\frac{1}{532}$ the importance of the model \rightarrow spectral trend comparisons, given the accuracy of the ⁵³³ AIRS observations.

⁵³⁴ The paper by (Raghuraman et al., 2023) shows similar figures, but in terms of spectral OLR trends encompassing the 0-2000 cm⁻¹ range, while (Huang et al., 2023) shows ⁵³⁶ similar plots for a slightly smaller time period (2002-2020) and using nadir L1B radiance ⁵³⁷ dataset which has no or minimal frequency corrections compared to the L1C set we use ⁵³⁸ in this paper. (Huang et al., 2023; Raghuraman et al., 2023) and our work all show, ei-⁵³⁹ ther in radiance or OLR space, (a) the increased observed radiance in the window channels, due to surface temperature increases (b) the \simeq -0.06 K \rightarrow yryr⁻¹ decrease in BT in $_{541}$ the 700-750 cm⁻¹ troposphere sounding region, which is due to the CO_2 amounts increas- $_{542}$ ing; we also see differences in the signs of the BT changes in the 650-700 cm⁻¹ strato- $_{543}$ spheric CO₂ and temperature channels for some latitudes between AIRS RT observations and ERA5 simulations (c) increases in the 1350-1640 cm⁻¹ water vapor sounding region seen in Figures 3 and 5, and (d) the 1280-1340 cm⁻¹ decreases are due to CH₄ ⁵⁴⁶ increases.

⁵⁴⁷ 4.1 Sample spectral closure comparisons using other monthly products

 Here we follow the earlier work of Huang et al. (2023) and convert the ERA5 monthly model fields to spectral radiances, after which we compute spectral trends for compar- ison to AIRS observations. Spectral closure calculations for the entire 20 year timeseries were also generated for the monthly MERRA2 model fields, as well as the monthly AIRS v7 L3 and CLIMCAPS L3 retrieved data products. Again only the monthly thermody- $\frac{5}{553}$ namics and surface temperature fields for all 72×64 tiles were used in these SARTA runs, with GHG changes added in for each timestep as described above. Spectral trends were then computed using Equation 2.

 We chose just one limited example here to illustrate the power of this approach for diagnosing which dataset is more accurate, given that the AIRS spectral trend accuracy is already established. Water vapor is highly variable in space and time, meaning wa- ter vapor retrievals using hyperspectral sounders radiances differ most from NWP fore- casts, in particular because of the typical ± 90 minute difference between observation and forecast, and is where these sounders typically provide the most information. Fig-⁵⁶² ure 5 show the globally averaged brightness temperature trends (in K $/\text{year}y^{-1}$) in the $_{563}$ 1350 - 1650 cm⁻¹ water vapor sounding region. The blue curve shows the trends from ₅₆₄ the AIRS observations used in this paper, while spectral trends constructed from the AIRS L3/ CLIMCAPS L3 retrievals are in red/yellow and the ERA5 model fields are in pur- ple. The AIRS observations and ERA5 constructed spectral trends are positive in this ₅₆₇ region, while the AIRS L3 and CLIMCAPS L3 trends are obviously different, being neg- ative in this water vapor sounding region. The subtle differences in these spectral trends ₅₆₉ arise from differences in the geophysical trends between observations and the models them- selves, and will be addressed in the following sections, where the retrieved and model sur- face temperature, and atmospheric temperature and water vapor geophysical trends will be compared and discussed.

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574 5 Testing the variability of representative points from NWP models reanalysis

Each sixteen day $3^{\circ} \times 5^{\circ}$ tile contains ~ 12000 observations, which means for each ⁵⁷⁶ tile about 600 daytime and 600 nightime observations are averaged to produce the Q0.90 $_{577}$ dataset per timestep. Conversely there are typically only ~ 240 monthly ERA5 0.25° points $_{578}$ per $3° \times 5°$ tile; for 1° resolution AIRS L3 and CLIMCAPS L3 there are even fewer (15) ₅₇₉ points per tile. This low number of points means we chose a simple solution of using the sso grid cell closest to the center of each $3° \times 5°$ tile for building the NWP and L3 geophys-⁵⁸¹ ical time series. This choice is validated below using the following test to see for exam-

Figure 5. Globally averaged spectral trends in the water vapor sounding region : AIRS L1C observations (blue) compared to spectral closure from the standard monthly AIRS L3 retrievals (red) and CLIMCAPS L3 (yellow) and from monthly ERA5 simulations (yellow). The reconstructed AIRS_RT trends very closely match the AIRS L1C observations and are not shown here.

⁵⁸² ple how surface temperature trends would be impacted as we changed the representa-⁵⁸³ tive point for the ERA5 model fields.

⁵⁸⁴ For the descending overpass we built complete sets of approximately 240 ERA5 points 585 per tile per month; at 0.25° resolution one of these is almost certainly at the tile center. ⁵⁸⁶ From these monthly sets, we could either directly read the tile center temperature (our ⁵⁸⁷ default), or compute the average surface temperature per tile, or compute the average ⁵⁸⁸ of the hottest 10% surface temperatures per tile. This was done for all 20 years (240 monthly ⁵⁸⁹ timesteps) after which the three timeseries were trended. Over ocean the differences between all three datasets as was typically $-0.001 \pm 0.005 \text{ K}$ /yeary $^{-1}$, while over land the $\text{differences were larger at about } 0.001 \pm 0.01 \text{ K } \frac{\text{year}}{\text{year}} \sim \text{m}^{\text{-1}}$. This is to be compared to $\sigma_{\rm{592}}$ mean trends of about 0.014 \pm 0.02 K $\frac{1}{\sqrt{y} \cdot y \cdot y \cdot \sqrt{y}}$ over ocean and 0.025 \pm 0.04 K $\frac{1}{\sqrt{y} \cdot y \cdot \sqrt{y}}$ $\frac{1}{2}$ land. In other words yr^{-1} over land : the spread of the ocean and land ERA5 surface ₅₉₄ temperature trends for the three methods, was about four times larger than the spread ₅₉₅ of the differences between the three methods. In what follows is much smaller than the sse mean trends. Given that there were far fewer re-analysis points in a grid box than tiled esity and the Coupled with the fact that choosing the 10% warmest profiles would provide an even smaller sample, we chose to use the tile center was thus chosen as to be ₅₉₉ the representative point to co-locate the model fields, when comparing against the tiled ⁶⁰⁰ observations.

⁶⁰¹ 6 Geophysical Trend Retrieval outline

⁶⁰² 6.1 Setting up the Retrieval Problem

 ϵ_{603} The observed spectral brightness temperature for a tile at any time t can be mod-⁶⁰⁴ eled as

$$
BT(\nu, t) = f(X(t), \epsilon(\nu, t), \theta(t)) + \text{NeDT}(\nu)
$$
\n(3)

605 where the state vector $X(t)$ has the following five geophysical state parameters : (1) sur- ϵ_{606} face temperature (ST), (2) atmospheric temperature profile T(z), (3) water vapor pro $\frac{607}{1000}$ file WV(z), (4) ozone profile O3(z) (5) greenhouse gas forcings (GHG) due to CO₂, CH₄ 608 and N₂O changing as a function of time t and $f(X(t), \epsilon, \theta, \nu)$ is the clear sky radiative ϵ_{09} transfer equation for channel center frequency ν . The spectral noise NeDT(ν) for a typ-⁶¹⁰ ical tropical "clear scene" is about 0.1 K in window region, increasing to about 1 K in ϵ_{01} the 15 μ m temperature sounding channels and about 0.2 K in the 6.7 μ m water vapor ⁶¹² sounding region, but the noise will vary as a function of scene temperature. We parameterize ϵ ₆₁₃ parametrize the GHGs using single numbers (such as ppm(t) for the CO₂ column), and 614 include the AIRS orbit and viewing angle geometry θ and the surface emissivity $ε(ν)$, ⁶¹⁵ while we omit forward model and spectroscopy errors. We ignore cloud scattering as well ⁶¹⁶ as the spatial variation of the state parameters, emissivity and scan angle geometry within ⁶¹⁷ a tile. Linearizing the above equation about the time averaged profile, the relationship ⁶¹⁸ between the observed spectral trends and desired thermodynamic trends is given by

$$
\frac{d\overline{BT(\nu)}}{dt} = \frac{\partial f}{\partial \overline{X}} \frac{d}{dt} \overline{X(t)} = K(\nu) \frac{d}{dt} \overline{X(t)} + K_{\text{emissivity}}(\overline{\nu}) \frac{d}{dt} \overbrace{e(t)}^0 \to K(\nu) \frac{d}{dt} \overline{X(t)} \tag{4}
$$

where the matrix $K(\nu)$ is the thermodynamic jacobian (surface temperature, air ϵ_{620} temperature and trace gases) and we ignore any orbit drifts (changes to θ), instrument ϵ_{621} changes (changes to $NeDT(\nu)$) and surface emissivity $(\epsilon(\nu))$; the last assumption is in- ϵ_{622} vestigated in a later section. The overbars on parameters X denotes this is a time av-⁶²³ erage (linear trend) that we are working with, and we have converted from radiances in ⁶²⁴ Equation 2 to brightness temperatures in Equations 3 and 4.

⁶²⁵ 6.2 Jacobian calculations

 F_{626} For a typical clear sky tropical sky atmosphere, the 800 - 1200 cm⁻¹ window re- ϵ_{627} gion has surface temperature (SKT) jacobians which are about $+0.5$ to $+0.75$ K per de- ϵ_{28} gree SKT change and -0.75 to -0.25 K per 10% change in column water vapor. The spec-⁶²⁹ tral variability in these window region jacobians is primarily due to reducing water con- ϵ_{500} tinuum absorption as you move from the 800 cm⁻¹ end to the 1200 cm⁻¹; consequently 631 the surface temperature jacobians becomes closer to unity and the column water jaco-⁶³² bians become closer to zero as water vapor amount decreases (drier atmospheres in the ⁶³³ mid-latitudes and polar regions). The hyperspectral channels used in this work help sep-⁶³⁴ arate out these two competing changes, which we validate against other datasets in this ⁶³⁵ study. As seen in Figure 4 typical magnitudes of the spectral trends on the left hand side ⁶³⁶ of Equation 4 are less than about 0.1 K per year. Equation 4 is in the usual inversion δ_{37} form $\delta y = K \delta x$, and the Optimal Estimation (Rodgers, 2000) solution used to solve ⁶³⁸ the anomaly time series in Strow et al. (2021) is also used here. The noise term used for ϵ_{639} the trend retrieval $NeDT(\nu)$ is not the instrument noise since each 16 day point in our ⁶⁴⁰ time series is averaged over hundreds of observations as earlier described; instead the un-⁶⁴¹ certainty is that due to inter-annual variability in the linear trends obtained from the ⁶⁴² trend fitting in Equation 2. Examples of typical noise values are shown in the bottom right hand panels panel of Figures 3and ??... 643

 ERA5 monthly model fields at tile centers, together with time varying concentra- ϵ_{45} tions of GHG such as CO_2 , were averaged over 20 years so jacobians could be computed. The GHG concentrations were a latitude dependent increase of about ∼ 2.2 ppmv/year $_{647}$ ppm yr⁻¹ for CO₂ derived from the CarbonTracker Peters et al. (2007) (CarbonTracker CT-NRT.v2023-4, http://carbontracker.noaa.gov) model data at 500 mb. Our pseudo- monochromatic line by line code kCARTA De Souza-Machado et al. (2018, 2020) was used with these averaged profiles to produce accurate analytic jacobians. The HITRAN 2020 line parameter database Gordon and Rothman (2022), together with MT-CKD 3.2 and $_{652}$ CO₂, CH₄ line mixing from the LBLRTM suite of models Clough et al. (2005) were used in the kCARTA optical depth database De Souza-Machado et al. (2018). A 12 month ge- ographical land-varying spectral emissivity database spanning one year from Zhou et al. (2011) was used, while ocean emissivity came from Masuda et al. (1988). The atmospheric temperature, water vapor and ozone profile jacobians, and the surface temperature and column jacobians for the GHG gases such as $CO₂$ and $CH₄$ and $N₂O$, were then convolved using the best estimate AIRS Spectral Response Functions Strow, Hannon, Weiler, et al. (2003).

 Tests done for this paper, together with the results in Strow et al. (2021), estab- lished that jacobians derived from MERRA2 versus ERA5 produced no significant dif- ferences in the context of retrieved trends or anomalies done for this paper, as the un- certainty in linear trends due to inter-annual variability dominates over any uncertainty (or differences between) model fields.

⁶⁶⁵ 6.3 Optimal Estimation Retrieval : State vector, covariance matrices ⁶⁶⁶ and a-priori

 $\frac{667}{1666}$ Using monthly ERA5 model fields averaged over 20 years, for each of the 64 \times 72 ⁶⁶⁸ tiles we computed analytic jacobians for the following (vector) atmospheric thermodynamic variables [fractional water vapor, fractional ozone and temperature] together with ϵ_{670} (scalar) surface temperature, where we retrieved fractional gas concentration trends $dfracX/dt =$ $\frac{671}{1/X_{avg}(z)dX_{avg}(z)}$ to keep all values in the state vector at about the same magni-⁶⁷² tude. A single iteration Optimal Estimation retrieval Rodgers (2000) is used to simul-⁶⁷³ taneously solve for the geophysical parameter trends. As in (Strow & DeSouza-Machado, 674 2020) the geophysical covariance uncertainty matrices are a combination of Tikonov and ⁶⁷⁵ covariance regularization. The uncertainties for the covariance matrices were typically [0.1,0.25,0.45] K /yr ✿✿✿✿ yr[−]¹ ⁶⁷⁶ for the surface/tropospheric/stratospheric temperature trends, σ_{677} and [0.04/0.02] $\sqrt{\text{year yr}^{-1}}$ for the fractional tropospheric/stratospheric water vapor trends. ⁶⁷⁸ Tikonov L1 regularization (Rodgers, 2000) also included, with the scalar factor multi- ϵ_{679} plying this regularization corresponding to about $1/10$ the covariance uncertainties. The ⁶⁸⁰ spectral uncertainties used in the retrievals come from the above mentioned trend un-⁶⁸¹ certainties. For completeness we note that a sequential retrieval (see for example Smith ⁶⁸² and Barnet (2020)) produces very similar geophysical trends.

 Here we emphasize four points about our geophysical trend retrievals, which sets $\frac{1}{2}$ us apart from trends derived from other datasets. Firstly the *a-priori* trend state vec- $\epsilon_{\rm 685}$ for is zero $\frac{\rm (dST/dt=dT(z)/dt=QQ(z)/dt=0)}{\rm (dST/dt=2C_z)}$ for all geophysical parameters, except for water vapor where we enforced constant (or slightly increasing) relative humidity as described below. This ensures traceability of our retrieval is straightforward especially wherever the AIRS instrument has sensitivity. For example the 300 - 800 mb water va- por trend retrievals will be based on the data only, thereby insulating us from any pos- ϵ_{000} sible *a-priori* information from *e.g.* climatology or NWP models, unlike the operational AIRS V7 or CLIMCAPS retrievals which use first guesses based on neural net and MERRA2 respectively.

 $\frac{693}{100}$ Secondly as seen in Figures 4 and 5, in the 15 μ m region there is a large spectral overlap signal $(-0.06 \text{ K } / \text{y} \text{y} \text{y}^{-1})$ from the increasing CO₂, which is much larger than the 695 expected atmospheric temperature trend $(0.01 \text{ K}/yr)$. The 20 year dataset contains inter-annual ⁶⁹⁶ variability whose noisy time series and correlations with for example temperature changes, $\frac{1}{\text{whibch-yr}^{-1}}$. These correlations makes it difficult to also retrieve these well mixed GHG. ⁶⁹⁸ Instead of attempting to solve for both GHG concentration changes and for temperature ese changes, we spectrally removed-jointly retrieve both temperatures changes and changes ro in well mixed GHGs such as CO_2 . We chose to focus on retrieving temperature changes ⁷⁰¹ only, by spectrally removing the effects of changing $edCO_2$, CH₄ and N₂O GHG con-⁷⁰² centrations, . This was done by using the GHG trends estimated from NOAA ESRL Car- $_{703}$ bonTracker data multiplied by the appropriate GHG gas column jacobian (CO₂,N₂O and $CH₄$ and CFC11, CFC12) computed as described above using the averaged over 20 years ⁷⁰⁵ ERA5 monthly profile for each tile.

 Thirdly instead of using all 100 layers described in the AIRS forward model Strow, Hannon, DeSouza-Machado, et al. (2003), we combine pairs of layers for a 50 atmospheric layer retrieval, as the AIRS radiances contain far fewer than 100 pieces of information $\frac{709}{200}$ (see e.g. Maddy and Barnet (2008); De Souza-Machado et al. (2018)).

 Fourthly, modern hyperspectral infrared sounders have highest sensitivity to tem- perature and water vapor in the mid-tropopause; see for example the averaging kernels in Irion et al. (2018). Using a zero fractional WV trends *a-priori* at all levels, it was fairly $_{713}$ straightforward to obtain fractional WV (z) trends close to those from the NWP model datasets in the $3000-850-300-850$ mb region. In order to improve our results in the low- est layers, we enforced a constant relative humidity approximation, which is a well-known, expected behavior under global climate change Soden and Held (2006); Sherwood et al. (2010). This was done by using the ignoring the contribution due to water vapor changes in the observed BT1231 trend, and using it as an approximation for air temperature trend over ocean; this allows us to compute an estimate of how the water vapor would need to change

$$
RH(T) = \frac{e}{e_{sat}(T)} \implies \delta(RH) = \frac{1}{e_{sat}(T)}\delta e - \frac{e}{e_{sat}^2(T)}\delta e_{sat}(T) = \frac{1}{e_{sat}(T)}\delta e - \frac{e}{e_{sat}(T)}\frac{L_v}{R_v}\frac{1}{T^2}\delta T
$$
(5)

where $e, e_{sat}(T)$ are the vapor pressures and we used $e_{sat}(T) = e_{s0}e^{\frac{L_v}{R_v}(\frac{1}{T_o}-\frac{1}{T})}$ (where L_v, R_v are latent heat of vaporization and gas constant respectively) to go from the ex-⁷²³ pression in the center to the expression on the right. If we expect the change in RH to be zero then $\frac{\delta e}{e} = \frac{L_v}{R_v} \frac{\delta T}{T^2}$, where we can use $\delta T/\delta t \sim d/dt B T 1231$. to approximate the 725 a-priori fractional vapor pressure rates (or a-priori fractional water vapor rates) between ⁷²⁶ surface and 850 mb, smoothly tailing to 0 in the upper atmosphere. Subsection 7.2 has ⁷²⁷ a similar discussion on a proposed method to alleviate the lack of sensitivity to upper $\frac{728}{128}$ atmosphere water vapor. Our default results in this paper are from using the MLS a -⁷²⁹ priori, unless otherwise stated.

⁷³⁰ 6.4 Testing on Synthetic Spectrasynthetic trend spectra made from ERA5 731 **Reanalysis monthly fields**

 We tested the retrieval code by using it on the simulated nighttime only ERA5 spec- tral trends, and compared to geophysical trends computed directly from the ERA5 re- analysis model. Spot checks of the spatial correlations of ERA5 fractional water vapor and temperature trends versus the trends retrieved from synthetic spectra/our retrieval algorithm, peaked at 500 mb with correlations of about 0.9, compared to 800 mb cor- relations of 0.80 and 0.55 for temperature and fractional water vapor trends respectively $\frac{738}{128}$ and 200 mb correlations of 0.89 and 0.69 for dT/dt, dWVfrac/dt. This is to be expected since a computation of the water vapor averaging kernels for infrared instruments for ar- bitrary atmospheric profiles typically shows they peak in the 300 mb - 850 mb range and decrease rapidly away from those regions; conversely the temperature averaging kernels stay relatively uniform through the free troposphere and above, though they also decrease close to the surface; see for example Irion et al. (2018); Smith and Barnet (2020); Wu et al. (2023).

⁷⁴⁵ Figure 6 shows a sample set of results using nightime ERA5 model output converted ⁷⁴⁶ to spectral trends as described above. The top panels (A) are always the atmospheric trends derived direct computed directly from the monthly ERA5 model fields, while the ⁷⁴⁸ bottom panels (B) are the atmospheric trends derived retrieved from the converted ERA5 ⁷⁴⁹ spectral brightness temperature trends. The left most panel is the atmospheric temper-⁷⁵⁰ ature trend comparison (both in K $/\overline{\text{yryr}}_{\text{av}}^{-1}$) while the rightmost panel is the fractional α_{751} atmospheric water vapor trend comparison (in $\sqrt{\text{yryr}_{\text{7}}^{-1}}$).

Figure 6. Comparing geophysical trends derived directly from ERA5 monthly nighttime fields (top) vs from the OEM retrieval applied to the spectral trends (bottom). Left panel is dT/dt (in K \rightarrow yearyr⁻¹</sub>) while rightmost panel is d(fracWV)/dt (colorbar in \rightarrow yr⁻¹_∞).

Table 1. Cosine weighted air temperature, skin temperature, fractional water vapor trends, together with uncertainties f. The "ERA5 direct" are directly from the ERA5 geophysical trends, while "ERA5 spectral" are trends retrieved from the converted ERA5 spectral trends.

⁷⁶¹ 6.5 Surface emissivity changes

⁷⁶² Equation 3 explicitly includes the surface emissivity in the equation of radiative ⁷⁶³ transfer; however Equation 4 assumes this is unchanging. Here we rewrite Equation 4 ⁷⁶⁴ as

$$
\frac{d\overline{BT(\nu)}}{dt} - K_{emissivity}(\nu)\frac{d}{dt}\overline{\epsilon(t)} \to \frac{d\overline{BT'(\nu)}}{dt} = K(\nu)\frac{d}{dt}\overline{X(t)}\tag{6}
$$

First we consider ocean emissivity changes Ocean emissivity has a dependence on ⁷⁶⁶ windspeed Masuda et al. (1988). Lin and Oey (2020) and other literature suggest wind speed increases of $+2.5 \text{ cm } / \text{s}/\text{year } \text{s}^{-1} \text{ yr}^{-1}$ have occured between 1993-2015 in the trop- τ_{68} ical Pacific, and smaller (or close to zero) values elsewhere. The monthly ERA5 $u10, v10$

⁷⁶⁹ 10 m speeds for the 20 year time period in this paper also showed the maximum abso- 770 lute trend was 0.09 m/s/year (over the Southern Ocean) while the global ocean mean and standard deviation were $0.006 \pm 0.022 \text{ m/s/years}^{-1} \text{yr}^{-1}$; The emissivity changes over ocean using a 0.025 m $\frac{1}{52}$ wind speed change are on average on the order of 1× 10^{-6} per year in the thermal infrared window (or about 0.0003 K $\sqrt{\text{yry}}_{\text{max}}^{-1}$ change in the ⁷⁷⁴ window region); assuming the optical properties of water do not substantially change with τ ⁷⁷⁵ the ∼ 0.02 K increases seen in all the datasets considered in this paper, these very small ⁷⁷⁶ emissivity changes are of no consequence.

⁷⁷⁷ Land emissivity changes were estimated as follows. A global monthly mean emis-⁷⁷⁸ sivity database, the Combined ASTER and MODIS Emissivity over Land (CAMEL v003) ⁷⁷⁹ has recently been released Borbas et al. (2018). We matched the tile centers to the database $\frac{780}{100}$ for the 20 \times 12 months spanning our 2002/09 - 2022/08 time period, and computed the emissivity trends over land; the results (not shown here) were on the order of -1×10^{-4} 781 ⁷⁸² and $+3\times10^{-4}$ in the 800-960 cm⁻¹ and 1100-1250 cm⁻¹ regions respectively, averaged ⁷⁸³ over the land observations. For each tile the $K_{emissivity}(\nu) \frac{d}{dt} \epsilon(t)$ term was estimated by running SARTA with the default emissivity, then differencing with the SARTA output ⁷⁸⁵ obtained when the emissivity trends were added on. Averaged over the planet, the spec-⁷⁸⁶ tral changes arising from these emissivity changes were much smaller than the spectral trends seen in Figure 3, about -0.001 K $\sqrt{\text{year} - y_x^{-1}}$ between 800-960 cm⁻¹ and about +0.002 ⁷⁸⁸ K \rightarrow $\frac{1}{2}$ $\frac{1}{2}$ on the 1100-1250 cm⁻¹ region (which we do not use in our retrieval, since ⁷⁸⁹ many of the channels are synthetic and the real channels are drifting (Strow et al., 2021)). $_{790}$ The land only results were roughly about three times these magnitudes. Using these emis- $\frac{791}{791}$ sivity jacobians on the left hand side of Equation 6 and running the retrieval on the ad- 792 justed spectral trends over land, resulted in about at most 0.01 K increases to the zon-⁷⁹³ ally averaged surface temperature changes over land; zonally averaged these largest differences were at about $40°N$ to $60°N$ and $-25°S$ to $+15°N$, due to emissivity decreases; $\text{the } 20^{\circ}$ N to $+35^{\circ}$ N region which included the Sahara and swathes of Asia, had emissiv-⁷⁹⁶ ity increases but the averaged-over-land temperature decreases were small, as there were ⁷⁹⁷ offsetting emissivity increases in other land areas at the same latitudes. We did not pur-⁷⁹⁸ sue the impact of these emissivity changes further as the CAMEL database is affected ⁷⁹⁹ by the stability of the MODIS data, and our results below will not include accounting ⁸⁰⁰ for changes in land emissivity.

⁸⁰¹ 7 Results

 The trends retrieved in the previous section using simulated radiance trends show that the retrieval package is working as expected. Here we apply our retrieval to observed AIRS L1C radiance trends and discuss the retrieved AIRS RT geophysical trends to those computed directly from the ERA5/MERRA2 model fields and AIRS L3/CLIMCAPS L3 products. We will have an expectation that since the simulated radiance trends had no noise added to them, the uncertainty in the spectral rates was lower than the actual ob- served spectral uncertainty; this will lead to larger uncertainties and/or errors in our re-trieval using observed radiance trends.

⁸¹⁰ We will make most comparisons against NWP models and L3 products in the con- B_{811} text of averages over the descending/night (N) and ascending/day (D) data since the MERRA2 $_{812}$ (and GISS) datasets are only available as a D/N average; the reader is referred to the ⁸¹³ Appendix where we show a few of the D-N differences. The results are shown in the or-⁸¹⁴ der of surface/column trends (surface temperature and column water), followed by zonal 815 averages of the atmospheric temperature and fractional water vapor trends.

816 7.1 Skin Temperature trends

⁸¹⁷ There are typically multiple (window) channels that are sensitive to a surface pres-⁸¹⁸ sure, meaning the radiances typically have more information content for the surface tem-

Figure 7. Surface temperature trends dSKT/dt averaged over day and night for AIRS RT, and from separately fitting the monthly data in ERA5, MERRA2, AIRS L3, CLIMCAPS L3 and GISS.

⁸¹⁹ perature (assuming the surface emissivity is well known and there are no clouds) rather $\frac{1}{820}$ than for example air temperature. Figure 7 shows the diurnally averaged day/night (D/N) ⁸²¹ surface temperature trends from 6 datasets : AIRS RT, AIRS L3, CLIMCAPS L3, ERA5, 822 MERRA2 and NASA GISTEMP. AIRS RT shows an overall global warming of $+0.021$ ⁸²³ K \rightarrow $\frac{1}{2}$ K \rightarrow $\frac{1}{2}$; the cooling trends include the tropical eastern Pacific and south of Green-⁸²⁴ land and tropical northern Atlantic. The rest of the datasets also show similar patterns ⁸²⁵ of cooling in the N. Atlantic Ocean, warming over the Arctic and some degree of cool-₈₂₆ ing over the Antarctic Ice Shelf/Southern Ocean as does AIRS RT. The AIRS v7 L3 827 shows some cooling over Central Africa and the Amazon not seen in the AIRS RT trends, where one could expect Deep Convective Clouds and possible cloud clearing issues. We 829 also point out the AIRS L3 product has many missing values of the western coasts of 830 N. and S. America, due to cloud clearing issues. MERRA2 shows more cooling over C. 831 Africa, and just like the AIRS v7 data, a lot of cooling near the Antarctic Ice Shelf. Of ⁸³² note here is that although CLIMCAPS uses MERRA2 as its first guess, their surface tem-⁸³³ perature trends are not similar, especially around the Antarctic where MERRA2 shows ⁸³⁴ strong cooling trends.Over the ocean GISS shows similar trends to what AIRS_RT trends ⁸³⁵ show. An earlier study of Land Surface Temperatures between 2003-2017 using MODIS ⁸³⁶ Prakash and Norouzi (2020) shows very similar large daytime cooling trends over parts ⁸³⁷ of central and western Indian subcontinent that we see from our retrieval as well as di-⁸³⁸ rectly from the BT1231 channel trends; for tiles that straddle both ocean and land the ⁸³⁹ quantile method picks up the hottest observations, which especially during summer are ⁸⁴⁰ mostly over the Indian subcontinent. For these reasons we also have confidence in our ⁸⁴¹ retrieved cooling trends over for example daytime continental Central/Eastern Africa, 842 which are different from the other four day/night datasets.

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⁸⁴⁴ The spatial correlations between AIRS \overline{RT} retrieved rates and the various datasets ⁸⁴⁵ is shown in Table 2 while the cosine weighted skin temperature trends are shown in Ta-⁸⁴⁶ ble 3. By adding in the uncertainty in the trends for any of the individual models or datasets, and then doing the cosine weighting, we estimate uncertainties of about \pm 0.015 K $/\text{yry}^{-1}$ 847 ⁸⁴⁸ for "ALL"; the uncertainties for "OCEAN" are typically about $2/3$ of that value, and for ⁸⁴⁹ "LAND" are about $4/3$ of that value. We emphasize here that we use all available NWP 850 and L3 model data when computing their trends for any grid box, while the AIRS RT ⁸⁵¹ uses only the hottest 10% of "clear" data; Strow and DeSouza-Machado (2020) showed ⁸⁵² that the tropical retrieved surface temperature trends and anomalies over ocean corre-⁸⁵³ lated very well with those from the ERA-I Sea Surface Temperature dataset.

| | | | ERA5 MERRA2 AIRSL3 CLIMCAPSL3 GISS | |
|------|------|------|------------------------------------|------|
| 0.72 | 0.59 | 0.80 | 0.89 | 0.77 |

Table 2. Correlations of average (nighttime,daytime) retrieved skin temperature trends from AIRS_RT, versus trends from models/products

| SKT trend K $\left\langle \frac{W}{\sqrt{2}}\right\rangle$ AIRS RT AIRS CLIMCAPS ERA5 MERRA2 | | | | | | GISS |
|--|-------|-------|-------|-------|----------|-------|
| ALL | 0.020 | 0.017 | 0.021 | 0.023 | 0.011 | 0.021 |
| TROPICS | 0.011 | 0.011 | 0.012 | 0.016 | 0.010 | 0.015 |
| MIDLATS | 0.029 | 0.020 | 0.028 | 0.026 | 0.020 | 0.026 |
| POLAR | 0.032 | 0.028 | 0.033 | 0.041 | -0.005 | 0.028 |
| OCEAN | 0.019 | 0.011 | 0.019 | 0.017 | 0.012 | 0.017 |
| LAND | 0.022 | 0.030 | 0.024 | 0.038 | 0.010 | 0.030 |

Table 3. Cosine weighted skin temperature trends; uncertainties are on the order of \pm 0.015 K as explained in the text.

 A notable outlier in this group is the MERRA2 trends, especially over land and ⁸⁵⁵ the Southern Ocean which are noticeable negative (blue) compared to the other datasets; ⁸⁵⁶ the agreement with tropical and mid-latitude oceans is much better. As noted earlier, the MERRA2 monthly trends come from a combination day/night dataset that was down- loaded, which as seen in Figure 7 consists of trends that are both positive and negative, combining to get a closer-to-zero global weighted trend. In addition MERRA2 is the only ⁸⁶⁰ one of the six that (a) does not have the extreme +0.15 K $/\sqrt{\text{year } \text{y}r^{-1}}$ warming in the northern polar region and (b) shows a lot of cooling in the Central African area. Using ERA5 monthly data, we devised a test similar to the one mentioned in Section 5 to de- termine if the differences between MERRA2 and ERA5 surface temperature trends could be due to the temporal sampling (once for MERRA2 versus eight times for ERA5). For each month we matched the eight ERA5 timesteps available per month to the tile cen- ters and then averaged the surface temperatures per month; the ensuing geophysical time- $\frac{867}{100}$ series was then trended. The day/night ERA5 average of Figure 7 was compared to these trends; of note are (a) we did not see the cooling in Africa and near the Antarctic that is seen in MERRA2 and (b) the main differences between the $1.30 \text{ am}/1.30 \text{ pm}$ average in the bottom middle (ERA5) panel were over land (all 5 continents); the histograms of t_{max} the differences showed the peak was typically close to 0 K \rightarrow $\frac{t}{\sqrt{2}}$, but the widths over ⁸⁷² land were about ± 0.02 K $\sqrt{y_{\text{F}}y_{\text{L}}-1}$ or less (compared to ± 0.005 K $\sqrt{y_{\text{F}}y_{\text{L}}-1}$ over ocean).

Figure 8. Zonally averaged surface temperature trends \underline{for} (left) sum of ocean and land point and (right) ocean only.

⁸⁷³ Both AIRS L3 and MERRA2 show cooling in the Southern Ocean; we note that although 874 MERRA2 is the *a-priori* for CLIMCAPS L3, their trends are different that those from 875 MERRA2; in fact AIRS RT shows the closest correlation to the observational CLIM-⁸⁷⁶ CAPS L3 trends. The AIRS L3 trends in the Southern Ocean region could arise because $\frac{877}{100}$ of problems identifying ice during the L2 retrieval (private communication : Evan Man-⁸⁷⁸ ning (JPL) and John Blaisdell (NASA GSFC)) though the MERRA2 trends also show ⁸⁷⁹ significant cooling in that region, where few surface observations from buoys poleward 880 of 60[°] exist to help resolve these differences (see for example Figure 10 in Haiden et al. 881 (2018) .

⁸⁸² Figure 8 shows the zonally averaged total (land+ocean) and ocean only surface tem-⁸⁸³ perature trends. Notice how the equator to midlatitude ocean trends are almost linear ⁸⁸⁴ for all datasets, with the slope for the northern hemisphere being about double that of the southern hemisphere (roughly 0.001 K \rightarrow $\frac{1}{2}$ per deg latitude). Again focusing on the right hand plot, the AIRS L3 trends are negative in the Southern Ocean regions, ess compared to the other 3 datasets, due to the cooling trends around the Antartic continent ass shown earlier, but then agrees with most of the other datasets over the Antartic; the MERRA2 trends significantly differ between -90 S and -50 S. MERRA2 and ERA5 also show slightly ese <u>smaller</u> warming trends in the Northern Polar, compared to the three AIRS-based datasets. 891

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 We point out that the trends seen in Figure 7 vary noticeably at more local, regional levels and furthermore this spatial variation can differ between daytime and nighttime, ⁸⁹⁵ evident in Figure A1 of Appendix Appendix A, and that the observational sets (AIRS RT, CLIMCAPS L3 and AIRS L3) had larger differences than ERA5. Discussing the pos-sible causes of this is outside the scope of the paper.

898 **7.2** <u>Addition of Microwave Limb Sounder</u> Water Vapor A-priori

899 The Microwave Limb Sounder (MLS), on board NASA's Aura platform, is designed of the atmosphere above 300 mb. We computed water vapor trends from on the L3 data produced for that instrument (above 300 mb) and used them as an a-prioritor the AIRS_RT retrieval. 902

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Figure 9. dWVfrac/dt (left) without and (right) with MLS *a-prioriin* the upper atmosphere

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Figure 9 shows the retrieved fractional water vapor trends when the *a-priori* trend ore in the upper atmosphere in the left and right panels were zero, or used MLS trends, respectively. es One sees that the additional information brought in by the instrument sensitive to upper or troposphere humidity, significantly changes the water vapor sounding especially in the es polar region by moving towards the MERRA2 and ERA5 fractional water vapor trends een in Figure 13. We note that the results shown in this paper use the MLS a-priori 910

⁹¹¹ 7.3 Column water vapor trends

⁹¹² Column water vapor trends provide an assessment of the water vapor retrieval qual-⁹¹³ ity in the lower atmosphere since this is dominated by the layers near the surface. For ⁹¹⁴ a hyperspectral infrared sounder over ocean the 1226 (Channel ID 1511) and 1231 (Channel 915 ID 1620) spectral points are similarly impacted by surface emissivity and absorption by the The water vapor information in the lowest layers is best retrieved using the weak water en lines in thermal infrared region. This part of the retrieval is significantly complicated by presence of nonzero surface temperature, air temperature and water equalar apport jacobians in this spectral region, meaning the AIRS instrument has much reduced ₉₂₀ sensitivity to the water vapor continuum. However the 1226 channel is on the wing of ⁹²¹ a weak water vapor line and has additional absorption from the atmospheric water vapor ⁹²² column. Subtracting the observed brightness temperatures of these two channels BT1231 923 - BT1226 is therefore a representative approximation to (but is not equal to) the column $\frac{924}{924}$ water, just as BT 1231 is a representative approximation to (but is not equal to) surface ₉₂₅ temperature. For example, using the simulated AIRS L1C clearsky radiance dataset over ⁹²⁶ ocean we constructed for this paper using ERA5 monthly fields, we can regress the ERA5 927 column water against the brightness temperature difference to obtain $mmw \sim 5.6$ (BT1231-BT1226) \leftarrow 1.0; over land the emissivity could vary rapidly enough that this approximation breaks down vapor amounts in these lowest layers. In addition the changing concentration of very ases minor gases such as CFC-11 and CFC-12 Strow and DeSouza-Machado (2020) are quite evident in the spectral trends, further complicating the water vapor trend retrieval for ⁹³² the lowest layers.

⁹³³ The left hand panel of

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⁹³⁵ Figure 10 shows the zonally averaged column water vapor trends, while the right ⁹³⁶ hand panel shows the zonally averaged BT1231 - BT1226 trend (notice the multiplication

Figure 10. Zonally averaged column water vapor trends for AIRS_RT, AIRS L3, CLIMCAPS L3, ERA5 and MERRA2.

 $\frac{1}{2}$ factor of 5.6 mentioned above will roughly equalize the y-axis of the two panels). The 938 gray curve is the AIRS L1C observations, while the black curve ₉₃₉ the retrieval; the rest of the curves come from the fast model simulations model/data fields. The expanded at the series of the error barsed in the series which are on the or-941 der of \pm 0.005 mm/year.

⁹⁴² The column water vapor trends for , AIRS L3, CLIMCAPS L3, ERA5 and MERRA2. ⁹⁴³ The left hand panel shows the zonal averages, while the right hand panel shows the BT1231-BT1226 944 zonally averaged trends.

⁹⁴⁵ Close examination of the right hand panel-AIRS_RT is from our retrievals while the rest are directly from the NWP or L3 model fields. Close examination shows the CLIM-947 CAPS L3 column water trend is nearly identical to the MERRA2 trend, as is also seen ⁹⁴⁸ in lower atmosphere water vapor trends shown later in Figure 13. Conversely the col-⁹⁴⁹ umn water vapor trends for AIRS L3 are negative in the lower troposphere in the mid-⁹⁵⁰ latitudes and tropics, which is not to be expected given that the surface temperature trends ⁹⁵¹ are positive. AIRS_RT nominally agrees with ERA5 and MERRA2 in the tropics and ₉₅₂ midlatitudes, but is smaller than either in the northern polar regions. <u>A reduced rate</u> oss for AIRS_RT is additionally seen in the 0-50 N latitudes, where there is a larger fraction of land (for which we do not use the assumption of constant relative humidity) compared to the Southern Hemisphere. Screening out the tiles over land slightly improves the agreement between reanalysis (ERA5, MERRA2) vs AIRS_RTcolumn water trends. Examination of the spectral trends in the window region does not shed any more insight into the differences, as the observation spectral trends and NWP reconstructed trends are very similar and ⁹⁵⁹ we are fitting the observed trends. The magnitudes and patterns look similar to the 2005-⁹⁶⁰ 2021 column water trends shown in Borger et al. (2022), which were derived using ob-⁹⁶¹ servations from the Ozone Monitoring Instrument (OMI). We point out their 16 year zon-⁹⁶² ally averaged trends look similar to the 20 year ERA5 zonally averaged column water ₉₆₃ trends between -60°S and -10°S, but become almost a factor of 2 larger between -10°S

Figure 11. The 400 mb fractional water vapor trends for (left) AIRS RT and (right) ERA5 show general agreement except in the Southern Polar Regions.

⁹⁶⁵ The column water trends are summarized in Table 4.

no MLSGLOBAL (cosine average) 0.029 TROPICAL 0.039

Table 4. Column water trends based on OMI data (16 years) and AIRS RT, ERA5 and MERRA2 (20 years). The units are in mm $\sqrt{\text{year}y}^{-1}$; the uncertainties are on the order of 0.1 mm $\frac{1}{\sqrt{2\pi}}$ for OMI and AIRS_RT, and half that for ERA5 and MERRA2, and AIRS L3 and CLIMCAPS L3.trends using MLS *a-priori* are shown in the table, as are trends without the MLS a-priori

966 D/N differences (not shown) for AIRS RT were on the order of \pm 0.005 mm $\frac{1}{\sqrt{year}}$ $\text{yr}^{-1}_{\text{max}}$ (with daytime trends being smaller over land), for AIRS L3 were on the order of ± 0.01 mm $\sqrt{\text{year}-\text{yr}}^{-1}$ or more (with larger values happening over the daytime tropi- γ_{969} cal oceans), while that for ERA5 and CLIMCAPS L3 were typically on the order of \pm ⁹⁷⁰ 0.03 mm $\sqrt{\text{year}-y}$ or less. Figure 11 shows the 400 mb fractional water vapor trends, 971 with the left panel being the AIRS_RT trends while the right panel is the ERA5 trends. ⁹⁷² Note that there is general agreement except in the Southern Polar region, which is as 973 also seen later in Figure 13 to some extent in the other two observational L3 datasets 974 (AIRS v3 and CLIMCAPS). This could be related to $\frac{a}{a}$ paper-work by Boisvert et al. (2019) ⁹⁷⁵ who showed decreasing evaporation from the Southern Ocean in the 2003-2016 period ⁹⁷⁶ due to increasing ice cover.

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⁹⁷⁸ 7.4 Zonal atmospheric temperature and water vapor trends

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Figure 12. Zonally averaged dT/dt shown in 5 panels. Horizontal axis is latitude while vertical axis is pressure. The y-limits are between 10 to 1000 mb, on a logarithmic scale.

 Figure 12 shows the zonally averaged atmospheric temperature trends from five of ⁹⁸¹ the datasets in Figures 7,10 above. In the troposphere the AIRS RT retrievals show the same general features as the trends from ERA5, though they begin to diverge in the strato- sphere and especially above that. In particular AIRS_RT does not show warming in the Southern Polar stratosphere; we have separately looked into seasonal trends and noted that our retrieved September/October/November temperature trends in the upper atmospheric Southern Polar regions are on the order of -0.12K /year✿✿✿✿ yr[−]¹ ⁹⁸⁶ , possibly lead- ing to an overall no net heating/cooling for the annual trends. In addition we point out that both our results and AIRS v7 L3 show a hint of cooling over the tropical surfaces. Note that CLIMCAPS is initialized by MERRA2, and their temperature trends are quite similar. AIRS v7 looks similar to AIRS RT except in the tropics where it almost has cooling in the lower troposphere and much more warming in the lower stratosphere. The correlations between AIRS_RT and the [AIRS L3, CLIMCAPS L3, MERRA2, ERA5] temperature trends of Figure 12 are [0.74,0.65,0.74,0.72] respectively.

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⁹⁹⁵ Figure 13 shows the zonally averaged atmospheric fractional water vapor trends $\frac{996}{1000}$ (d/dt WV(z,t)/ \langle WV(z,t)>). The five panels are markedly different from one another. ⁹⁹⁷ The AIRS RT trends resemble those of ERA5 in the tropical troposphere, though we ⁹⁹⁸ do not have drying in the lower tropical layers. Conversely, the observed trends in the Southern Polar (AIRS L3, CLIMCAPS L3 and AIRS RT) show drying rather than wet-₁₀₀₀ ting, though AIRS RT is less than that of CLIMCAPS/MERRA2. AIRS RT is an out-¹⁰⁰¹ lier in the upper polar atmosphere trends, as both the signals and the jacobians are close ¹⁰⁰² to zero. Of some concern is a little bit of drying in the northern polar region, where there

Figure 13. Zonally averaged dWV frac/dt shown in 5 panels. Horizontal axis is latitude while vertical axis is pressure. The y -limits are between 100 to 1000 mb, on a linear scale.

 are low H₂O amounts leading to small jacobians. CLIMCAPS v2 looks quite similar to the MERRA2 trends. AIRSv7 shows substantial drying in the lower troposphere, and considerable wetting in the upper troposphere, compared to any of the other datasets. 1006 Spectral closure studies (using the AIRS v7 H_2O trend \times the H_2O jacobians derived above from ERA5 average profiles) are not shown here, but differ noticeably from the CCR trends f_{1008} from AIRS v7 in the 1300-1600 cm⁻¹ region, indicating there are inadequacies in the AIRS V7 water vapor retrievals. The correlations between AIRS_RT and the [AIRS L3, CLIM- CAPS L3, MERRA2, ERA5] fractional water vapor trends of Figure 13 (limited to 100 mb, 1000 mb) are [0.65,0.24,0.36,0.58] respectively.

7.5 Addition of Microwave Limb Sounder Water Vapor A-priori

 The Microwave Limb Sounder (MLS), on board NASA's Aura platform, is designed for sounding of the atmosphere above 300 mb. We computed water vapor trends from the L3 data produced for that instrument (above 300 mb) and used them as an for the retrieval.

dWVfrac/dt (left) without and (right) with MLS in the upper atmosphere

 Figure 9 shows the retrieved fractional water vapor trends when the trend in the 1019 upper atmosphere in the left and right panels were zero, or used MLS trends, respectively. ¹⁰²⁰ One sees that the additional information brought in by the instrument sensitive to upper troposphere humidity, significantly changes the water vapor sounding especially in the polar region by moving towards the MERRA2 and ERA5 fractional water vapor trends seen in Figure 13.

8 Uncertainty

 The uncertainties for the AIRS v7 geophysical products are impacted by radiance noise amplification due to cloud clearing Susskind et al. (2003) and the neural net first guess, while state vector errors are estimated based on regressions. CLIMCAPS L2 geo- physical products are similarly impacted by cloud clearing noise in the radiances, but these are fully propagated together with geophysical error estimates from the MERRA2 first guess, through the retrieval algorithm which uses Optimal Estimation Smith and 1031 Barnet (2020). No estimate of uncertainties are available for the monthly L3 products.

 The uncertainties for the AIRS_RT trends is much more straightforward : the spec- tral uncertainties shown in Figure 4 are used together with the state vector covariance matrices to generate the uncertainty matrix using the relevant equations of Optimal Es- timation (Rodgers, 2000); we use the diagonal elements for the final uncertainties. Pan- $_{1037}$ els (A) and (C) of Figure 14 shows the zonally averaged (D/N) uncertainties as a func- tion of pressure and latitude. Inspection of the radiance trends uncertainties shown in the center panel of Figure 4 shows the upper atmosphere temperature sounding region $_{1040}$ (650-700 cm⁻¹) has much larger uncertainty in the polar regions. The instrument and spectroscopy characteristics, coupled with these observational uncertainties, are such that for temperature the smallest errors are in the tropics while the largest errors are in po- lar upper atmosphere, which are the regions below 100 mb where the ERA5 trends differ most from AIRS RT trends. Similarly for water vapor the larger errors are in the lower atmosphere and above about 300 mb; the constant RH assumption and MLS a-priori help alleviate the errors.

 $T₁₀₄₇$ The $h = ztest(trend, \mu = 0, trend \ uncertainty)Z-test$ confirmed this picture, as seen $\frac{1}{4!}$ in panels (B) and (D) of Figure $\frac{4!}{4!}$, which show the temperature and fractional water vapor trends, together with black dots marking the (latitude,altitude) points where the ¹⁰⁵⁰ **zero trend null hypothesis at the default significance level of** trends are larger than the iosi uncertainty in the trends, at the 5% was rejected significance level. This happens in panel (B) for the temperature trends in most of the tropical/mid-latitude free troposphere (and stratosphere) but not at the southern polar stratosphere; and in panel (D) for fractional water vapor trends in the 200-600 mb range, from the Southern Polar region to about $_{1055}$ +60 N latitude, and some spots in the Northern Polar.

9 Discussion

 In general for surface temperature trends, the disagreements between the six sets shown in Figure 7 are over the polar regions and over land (especially over the Amazon and Central Africa) and are smallest over tropical and mid-latitude oceans, indicating the best agreements, except for slightly larger differences off the western coast of the Amer- icas and Africa (which have a prevalence of MBL clouds). The atmospheric temperature trends in general agreed except for the upper atmosphere polar regions and in the high altitudes (less than about 200 mb). Similarly fractional water vapor trends differed most in the upper atmosphere (200 mb and above) and in the tropical/mid-latitude 600-800 mb region. A quick glance at Figure 13 shows the former is due to lower sensitivity to upper atmosphere water vapor, leading the AIRS_RT retrievals to have low values while the AIRS L2 retrieval is initialized by a neural net; conversely the latter is due to the AIRS L3 retrieval being negative while the rest were mainly positive. Similarly the AIRS_RT retrieval differs above the Antarctic continent.

 In general the observed surface temperature trends from the AIRS_RT retrievals agree with the ERA5 and MERRA2 trends, as well as the NASA GISS trends, except

Figure 14. Zonally averaged D/N plots of (A) temperature uncertainties in K $+{\rm yr}^{-1}_{\sim}\$ and (B) temperature trends in K $\rightarrow yr^{-1}$ together with null hypothesis. (C) and (D) are the same except for fractional water vapor uncertainty and trends in 1/year. See text for more detailed explanation.

 in the Southern Antarctic. That is a region where there are few surface observations; for retrievals there are competing effects of using ice vs ocean surface emissivity. Overall, the AIRS_RT retrieved surface temperature trends are typically in between ERA5 and $MERRA2$ for land + ocean in all regimes (tropical, midlatitude and polar), though slightly larger overall for ocean than the two reanalysis datasets; in general they are closer to the 1078 ERA5 trends than the MERRA2 trends.

 Strow et al. (2021) demonstrated that the long- and medium- wave channels of the 1080 AIRS instrument are radiometrically stable to better than 0.002-0.003 K \rightarrow yeary \sim , which is much smaller than the surface and tropospheric temperature trends in the reanaly-1082 sis models, AIRS L3 data and our retrieved trends. After A separate analysis of spectral trend uncertainties after 05,10,15,20 years of observations, Figure ?? shows the trend spectral uncertainties-years (not shown here) show that these uncertainties have been steadily decreasing and are now approaching this number, as can be seen in the bottom left panel 1086 of Figure 3. Furthermore, though we cannot guarantee only cloud free scenes in our cho- sen Q0.90 dataset used in this paper, the high correlations between other dataset surface trends compared to ours, is a good indication that our results come from mostly cloud-free scenes, or scenes whose clouds have negligible impact on our results.

 The observed zonal temperature trends agree with those from the models and the AIRS L3 products, except in the polar regions. Again this could be an issue of using slightly incorrect surface emissivity for the AIRS_RT retrievals. In addition we point out that since there is very little water vapor, the temperature jacobians near the surface are quite small in magnitude (compared to more humid atmospheres) and so it is difficult to sep- arate out the effects of surface temperature trends versus lower atmosphere temperature and H2O trends. The quantile construction used in this paper means that for example tiles straddling the subcontinent of India and the ocean will preferentially pick the land surface observations for daytime, which could lead to misleading trends on these coastal tiles. It is possible to subdivide the $3° \times 5°$ tiles into for example $1° \times 1°$ grids and do the analysis, but the number of observations per small grid cell would drop, leading to more noise in the retrieved trend.

 The AIRS_RT retrieved absolute column water trends are equal to/slightly larger than ERA5/MERRA2 in the tropics and below both of them in the midlatitudes; AIRS_RT ocean column water trends were slightly smaller than both ERA5 and MERRA2 over ocean, and in-between them over land. We note the difficulties we have retrieving H_2O close to the surface and in the upper atmosphere. This is simply a consequence of the sensitivity of the infrared sounder, namely most of the averaging kernels peak in the 300- 600 mb range. AIRS_RT column water trends agree with those from ERA5 and MERRA2 column water trends in the tropics; nevertheless even with expected lowered sensitivity to water vapor in the lower altitudes, we were able to retrieve similar column water va- por trends to the NWP models both in the tropics and in the mid-latitudes. The dif- ferences become more acute in the polar regions since the low average amounts of wa- ter vapor mean the water vapor jacobians are very small, as were the observed trends in the WV channels. However, we point out that our column water trends, which are both quite sensitive to water vapor in the lower atmosphere, are in good agreement with those 1116 from NWP models.

 We point our here that our results are relatively robust to changes in the covari- ance or Tikonov parameter settings. For instance changing them by factors of two would ₁₁₁₉ keep the trends about the same, though of course the uncertainties would change.

 Given the complex numerical algorithms used in both the reanalysis models and the AIRS L3 retrievals as well as those in the AIRS_RT trends, it is difficult to offer pre- cise explanations for any of the trends shown above. There are however a few general points that can be made. The first is that since infrared instruments are sensitive to the 300-800 mb region and lose sensitivity outside this, the retrievals from AIRS_RT and

1125 AIRS L3 have difficulties with water vapor in the lower (Planetary Boundary Layer) and upper troposphere *(lower stratosphere*. One way to mitigate this is to use trended data from external sources ; for in the a-priori, while keeping the a-prioritrends for all other $_{1128}$ parameters as 0. For example we have shown we can use the MLS data above 300 mb without significantly degrading the AIRS_RT retrieval in the middle and lower atmo- sphere; conversely the CLIMCAPS retrievals are initialized by MERRA2 and while they can pull out weather signals, their L3 trends are still quite closely tied to the MERRA2 trends. The tropical and mid-latitude ocean surface temperature trends from the numer- ical models that assimilate data, L3 products and AIRS_RT are very similar; however $_{1134}$ they start to show differences where there are few in-situ data combined with problems with ice identification (surface emissivity)/cold temperatures which exacerbate the drift-ing AIRS detector problems Strow et al. (2021), such as the Arctic and Southern Ocean.

1137 10 Conclusions

 We have described a novel method to obtain global thermodynamic atmospheric climate trends, starting from infrared allsky hyperspectral observations which are then subset for "nominally clear" scenes. Our retrieved trends are derived using trends from well characterized (radiometrically stable) radiances and from zero a-priori (except for a constant relative humidity assumption). This makes them much more direct and trace- able than trends from traditional L2 retrieval algorithms, which use complicated *a-priori* information. We also did "radiative closure" tests by running the monthly NWP or L3 fields through a radiative transfer model to compare the spectral trends so obtained against the observed spectral trends, which showed the most disagreement in the water vapor sounding regions.

 The temperature and water vapor trends retrieved from the "nominally clear" ra- diance trends resemble those computed from monthly ERA5 and MERRA2 reanalysis. The radiative spectral closure helps identify the cause of differences in the geophysical trends, rather than solely attributing them to deficiencies (eg the well known reduced sensitivity to water vapor near the boundary layer and above 200 mb) with our retrieval. For example the AIRS_RT temperature trends are quite similar to the reanalysis (MERRA2/ERA5) trends, while the water vapor (and/or Relative Humidity) trends are quite different, es- pecially in the lower troposphere and upper troposphere, which is clearly manifest as dif-ferences in the spectral trends in the water vapor sounding region.

¹¹⁵⁷ The 20 years of AIRS observations were binned into nominal 3×5 degree grid boxes covering the planet, with a time step of 16 days, from which anomalies and trends were obtained. To alleviate the reduced sensitivity of hyperspectral sounders to water vapor in the lower atmosphere we used an assumption of 0.01 increase in relative humidity to initialize the a-priori lower atmosphere fractional water vapor rates, while we similarly used Microwave Limb Sounder trends as an a-priori to address the high altitude water vapor deficiencies caused by lower sensitivity to upper atmosphere water vapor. New or updated time dependent surface emissivity databases may become available in the fu- ture, enabling us to include those effects into Equation 4. Problems in the polar regions and Planetary Boundary Layer water vapor retrievals will be harder to overcome since there is very little sensitivity to water vapor in these regions, together with fewer obser-vations to compare against, though more work is planned to address both of these.

 In this paper we used the 90th quantile (Q0.90) nominally "hottest" observed BT1231 data to form a time series over which to obtain radiance trends, after establishing that the spectral trends from this quantile differed by less than about \pm 0.0015 K $/\text{yryr}^{-1}$ from the 50th (or average) quantile. In the future we plan to base the data subset se- lection on MODIS cloud products (obtained at 1 km resolution compared to the AIRS 15 km resolution). In any case the AIRS L1C Q0.90 spectral trends used for the AIRS_RT results are very comparable to trends from quality assured binned AIRS CCR data Manning (2022). The quantile method allows us to select which data to use in the trends : we have explored doing the trend retrievals using the cloud fields contained in ERA5, together with the TwoSlab cloud algorithm De Souza-Machado et al. (2018) to compute jacobians when clouds are present, together with trends from the Q0.50 dataset described above. The retrieved geophysical trends resemble those described above in the mid to upper at- mosphere, and differ in the lower atmosphere, but more work is needed on this and is not discussed further. Longwave clear sky flux trends (both outgoing top-of-atmosphere and incoming bottom-of-atmosphere) and climate feedbacks will be discussed in a sep-arate paper.

 While the Aqua platform is scheduled to be terminated within the next few years, copies of near identical CrIS instruments are already in orbit, and more will be launched over the next few years, till at least 2040. The Climate Hyperspectral Infrared Radiance Product (CHIRP) Strow et al. (2021) will seamlessly combine the AIRS data between 2002-2015 to CrIS data from 2015-2040 to obtain a 40 year observational radiance record over which to study climate. This availability means that AIRS_RT and future AIRS/CrIS versions, is well positioned to enable climate analysis of geophysical trends for years to come.

Appendix A Day versus Night surface temperature trend differences

 Figure A1 shows the (top) daytime and (middle) nighttime surface temperature trends; from left to right the datasets are (observational) AIRS_RT, AIRS L3, CLIM- CAPS L3 and (reanalysis) ERA5. In general the AIRS observational datasets show en- hanced daytime cooling over the Indian subcontinent and Central Africa, compared to the ERA5 model; they also show daytime warming trends over continental Europe and central Asia and the Amazon are larger than during the nightime. With the large ocean heat capacity and smaller land heat capacity, the land is expected to show more of a di- urnal cycle than ocean. ERA5 sees warming over Eastern/Central Africa during daytime while the observational datasets see cooling. Similarly the three observational datasets see more daytime cooling over the Indian sub-continent and south eastern Australia than does ERA5; we omit more detailed analysis in this paper. During the nighttime, the AIRS L3 product has cooling over C. Africa and parts of the Amazon. The day-night differences are seen in the bottom row of the same figure. Note the colorbar is the same for ₁₂₀₇ all three rows. The differences are close to zero over the ocean. AIRS RT and CLIM- CAPS L3 see more daytime cooling over E. Africa and the Indian subcontinent. Over- all the magnitude of the day - night differences for the observations are larger for the AIRS observational datasets than for ERA5. ERA5 also sees negative differences over Central Asia compared to the AIRS observational datasets, which see positive differences (higher surface temperature trends during the daytime).

 The atmospheric temperature and fractional water vapor day-night differences are quite small (compared to the average values) and not shown here; AIRS L3 shows no- ticeable more wetting of the 600-800 mb region during daytime versus nightime, com-pared to the other three.

¹²¹⁸ Data availabilityQpen Research Section

 The AIRS L3 and CLIMCAPS L3 data products, as well as the AIRS L1C radi- ances are freely available to the public on the NASA servers. MERRA2 and ERA5 and GISTEMP model output are also freely available.

Figure A1. Top two rows : The (top) day and (middle) night surface temperature trends for AIRS_RT, AIRS L3, CLIMCAPS L3 and ERA5. Third row (bottom) is the D-N difference.

1222 **Author contribution**

1228 **Competing interests**

1229 The authors declare that they have no conflict of interest.

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